

PROBLEM SET 5 Due Friday October 28
Physics 115B- FALL 2011

Points
Assigned

This will be the last problem set before the Midterm Exam on Wednesday November 2.

Analytic:

- 10 [1.] Griffiths Problem 4.18
- 10 [2.] Griffiths Problem 4.19
- 10 [3.] Griffiths Problem 4.22
- 10 [4.] Griffiths Problem 4.24
- 10 [5.] Griffiths Problem 4.27

Comment: For the first part of the course, as you develop skill in programming, the computational problems will not necessarily have anything to do with quantum mechanics.

[6.] Write a C or C++ program to solve the discretized diffusion equation. One approach is to use two one dimensional arrays. First have a loop which gets the new density rhonew,

$$\text{rhonew}[n] = \text{rho}[n] + \frac{Ddt}{dx^2} (\text{rho}[n+1] - 2\text{rho}[n] + \text{rho}[n-1])$$

Then follow it with a loop which resets the density to these newly computed values.

$$\text{rho}[n] = \text{rhonew}[n]$$

These loops (over n) need themselves to be enclosed by a loop which iterates over some number of time steps: "nested loops". Alternatively you could use a two dimensional array,

$$\text{rhonew}[n][m+1] = \text{rho}[n][m] + \frac{Ddt}{dx^2} (\text{rho}[n+1][m] - 2\text{rho}[n][m] - \text{rho}[n-1][m])$$

(Here m is the time index.) What are the advantages/disadvantages of the two approaches?

[7.] So far in class we have not discussed the feasibility of numerical calculations. How can you estimate whether a calculation will be doable? A good crude approach is to think that the clock speed of the chip in your computer (about one GHz) means that the compute can do 10^9 operations (addition, subtraction, ...) per second. How long should your Kepler problem have taken to execute? Suppose you want to simulate the motion of 10^{11} stars in a galaxy with molecular dynamics. Could your computer do that problem? (How many time steps do you want to do?) If you can't simulate 10^{11} stars, how many could you study? (How long are you willing to run your computer? Will you use a super-computer with 10^5 processors? A petascale computer is described here: http://en.wikipedia.org/wiki/Blue_Waters).

70 TOTAL

#1 Griffiths Problem 4.18

L_+ is hermitian conjugate of L_- :

$$\begin{aligned}\langle \psi | L_+ \psi \rangle &= \langle \psi | (L_x + iL_y) \psi \rangle = \langle \psi | L_x | \psi \rangle + i \langle \psi | L_y | \psi \rangle \\ &= \langle (L_x - iL_y) \psi | \psi \rangle = \langle L_- \psi | \psi \rangle\end{aligned}$$

So L_+ and L_- are hermitian conjugates.

$$4.112 \Rightarrow L^2 = L_+ L_- + L_z^2 - \hbar L_z \quad \boxed{L_{\pm} f_l^m = (A_l^m) f_l^m}$$

For $L_+ L_-$:

$$\langle f_l^m | L_+ L_- | f_l^m \rangle = (A_l^m)^2 \langle f_l^{m+1} | f_l^{m+1} \rangle = (A_l^m)^2$$

$$\begin{aligned}\langle f_l^m | L^2 - L_z^2 + \hbar L_z | f_l^m \rangle \\ = (\hbar^2 l(l+1) - m^2 \hbar^2 + m \hbar^2) \langle f_l^m | f_l^m \rangle\end{aligned}$$

$$\text{So } \boxed{(A_l^m)^2 = \hbar^2 [l(l+1) - m(m-1)]^{1/2} \text{ for } L_+ L_-}$$

For $L_- L_+$:

$$\begin{aligned}\langle f_l^m | L_- L_+ | f_l^m \rangle &= \langle f_l^m | L^2 - L_z^2 - \hbar L_z | f_l^m \rangle \\ &= \hbar^2 l(l+1) - m^2 \hbar^2 - m \hbar^2 \langle f_l^m | f_l^m \rangle\end{aligned}$$

$$\Rightarrow \boxed{(A_l^m)^2 = \hbar^2 [l(l+1) - m(m+1)]^{1/2} \text{ for } L_- L_+}$$

#2 Griffiths Problem 4.19

$$L_x = y p_z - z p_y \quad L_y = z p_x - x p_z \quad L_z = x p_y - y p_x$$

$$\textcircled{a} \quad [L_z, x] = [x p_y - y p_x, x] = [x p_y, x] - [y p_x, x] = -y [p_x, x] = -y(-i\hbar) = \boxed{i\hbar y}$$

$$[L_z, y] = [x p_y - y p_x, y] = [x p_y, y] - [y p_x, y] = x [p_y, y] = x(-i\hbar) = \boxed{-i\hbar x}$$

$$[L_z, z] = [x p_y - y p_x, z] = [x p_y, z] - [y p_x, z] = \boxed{0}$$

$$[L_z, p_x] = [x p_y - y p_x, p_x] = [x p_y, p_x] - [y p_x, p_x] = [x, p_x] p_y = \boxed{i\hbar p_y}$$

$$[L_z, p_y] = [x p_y - y p_x, p_y] = [x p_y, p_y] - [y p_x, p_y] = -[y, p_y] p_x = \boxed{-i\hbar p_x}$$

$$[L_z, p_z] = [x p_y - y p_x, p_z] = [x p_y, p_z] - [y p_x, p_z] = \boxed{0}$$

$$\begin{aligned} \textcircled{b} \quad [L_z, L_x] &= [L_z, y p_z - z p_y] = [L_z, y p_z] - [L_z, z p_y] \\ &= y [L_z, p_z] + [L_z, y] p_z - z [L_z, p_y] - [L_z, z] p_y \\ &= 0 + i\hbar x p_z - z(-i\hbar p_x) + 0 = i\hbar(z p_x - x p_z) = \boxed{i\hbar L_y} \end{aligned}$$

$$\begin{aligned} \textcircled{c} \quad [L_z, \mathbf{r}^2] &= [L_z, x^2] + [L_z, y^2] + [L_z, z^2] \\ &= x [L_z, x] + [L_z, x] x + y [L_z, y] + [L_z, y] y + z [L_z, z] + [L_z, z] z \\ &= x(i\hbar y) + (i\hbar y)x + y(-i\hbar x) + (-i\hbar x)y + 0 + 0 = \boxed{0} \\ [L_z, \mathbf{p}^2] &= [L_z, p_x^2] + [L_z, p_y^2] + [L_z, p_z^2] \\ &= p_x [L_z, p_x] + [L_z, p_x] p_x + p_y [L_z, p_y] + [L_z, p_y] p_y + p_z [L_z, p_z] + [L_z, p_z] p_z \\ &= p_x(i\hbar p_y) + (i\hbar p_y) p_x + p_y(-i\hbar p_x) + (-i\hbar p_x) p_y + 0 = \boxed{0} \end{aligned}$$

\textcircled{d} L_x, L_y, L_z commute with \mathbf{p}^2 from (c) above
 L_x, L_y, L_z commute with \mathbf{r}^2 (from above) and thus commute with $V(r)$ (a function of \mathbf{r}^2).

#3 Griffiths Problem 4.22

(a) $L_+ Y_l^l = 0$ since $m=l$ is largest m value possible & L_+ destroys state.

(b) $L_z Y_l^l = \hbar l Y_l^l \Rightarrow -i\hbar \frac{\partial}{\partial \phi} Y_l^l = \hbar l Y_l^l$

Let $Y_l^l = A \Theta(\theta) \Phi(\phi)$

So $\frac{1}{\Phi} \frac{\partial \Phi}{\partial \phi} = il \Rightarrow \Phi(\phi) = e^{il\phi}$ (Φ part)

Eqn 130: $L_+ Y_l^l = \hbar e^{i\phi} \left(\frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) Y_l^l = 0$ (from (a))

$Y_l^l = A \Theta(\theta) e^{il\phi} \Rightarrow \left(\frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) A \Theta(\theta) e^{il\phi} =$

$A e^{il\phi} \frac{\partial \Theta}{\partial \theta} + i A \Theta \cot \theta \frac{\partial}{\partial \phi} e^{il\phi} = 0$

$\Rightarrow A e^{il\phi} \frac{\partial \Theta}{\partial \theta} - A \Theta e^{il\phi} (l) \cot \theta = 0$

$\Rightarrow \frac{\partial \Theta}{\partial \theta} - \Theta l \cot \theta = 0$

$\frac{\partial \Theta}{\Theta} = l \cot \theta \Rightarrow \ln(\Theta) = \int l \frac{\cos \theta}{\sin \theta} d\theta = l \ln(\sin \theta)$

$\Rightarrow \Theta = (\sin \theta)^l$

So, $Y_l^l = A \Theta(\theta) \Phi(\phi) = A (\sin \theta)^l e^{il\phi} = A (\sin \theta e^{i\phi})^l$

(c) $\langle Y_l^l | Y_l^l \rangle = \int_0^{2\pi} e^{il\phi} e^{-il\phi} d\phi \int_0^\pi A^2 \sin^l \theta \sin \theta d\theta = 2\pi A^2 \int_0^\pi \sin^{2l+1} \theta d\theta$

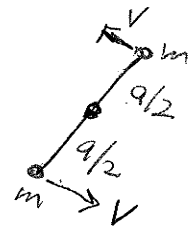
$= 2\pi A^2 \frac{(2)(2l)!!}{(2l+1)!!} = 4\pi A^2 \frac{2 \cdot 4 \cdot 6 \cdots 2l}{1 \cdot 3 \cdot 5 \cdots (2l+1)} = 4\pi A^2 \frac{(2 \cdot 4 \cdot 6 \cdots 2l)}{1 \cdot 3 \cdot 5 \cdots (2l+1)} \times \frac{(2 \cdot 4 \cdot 6 \cdots 2l)}{(2 \cdot 4 \cdot 6 \cdots 2l)}$

$= 4\pi A^2 \frac{(2^l l!)^2}{(2l+1)!} \Rightarrow A = \frac{1}{2^{l+1} l!} \sqrt{\frac{(2l+1)!}{\pi \cdot 1!}} \quad \checkmark \quad (3)$

#4

Griffiths Problem 4.24 $r \times p$

$$\textcircled{a} \quad L = r \times mv = a m v$$



Hamiltonian (Energy):

$$H = KE = \frac{1}{2} m v^2 + \frac{1}{2} m v^2 = m v^2$$

$$\Rightarrow \boxed{H = \frac{L^2}{m a^2}}$$

Eigenvalues of L^2 are $\hbar^2 l(l+1)$

$$H\psi = E\psi \Rightarrow \boxed{E_l = \frac{\hbar^2 l(l+1)}{m a^2}} \quad l = 0, 1, 2, \dots$$

⑥ Eigenfunctions are those of the L^2 operator; spherical harmonics:

$$\boxed{|\Psi_{lm}\rangle = Y_l^m(\theta, \phi)} \quad \text{with} \quad \boxed{2l+1}$$

degenerate energies

#5 Griffiths Problem 4.27

$$\chi = A \begin{pmatrix} 3i \\ 4 \end{pmatrix}$$

$$(a) \langle \chi | \chi \rangle = A^2 (-3i \ 4) \begin{pmatrix} 3i \\ 4 \end{pmatrix} = A^2 25 = 1$$

$$A = \frac{1}{5}$$

$$(b) \langle S_x \rangle = \chi^\dagger \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \chi = \frac{1}{25} [(-3i \ 4) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 3i \\ 4 \end{pmatrix}] = \underline{\underline{0}}$$

$$\langle S_y \rangle = \frac{1}{25} (-3i \ 4) \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 3i \\ 4 \end{pmatrix} = \frac{1}{25} (-3i \ 4) \begin{pmatrix} -4i \\ -3 \end{pmatrix} = \frac{\hbar}{50} \begin{pmatrix} -24 \\ -12 \end{pmatrix} = \underline{\underline{\frac{-12\hbar}{25}}}$$

$$\langle S_z \rangle = \frac{1}{25} (-3i \ 4) \left(\frac{\hbar}{2}\right) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 3i \\ 4 \end{pmatrix} = \frac{\hbar}{50} (-3i \ 4) \begin{pmatrix} 3i \\ -4 \end{pmatrix} = \underline{\underline{\frac{-7\hbar}{50}}}$$

(c) For spin $\frac{1}{2}$ $\langle S_j^2 \rangle = \chi^\dagger \left(\frac{\hbar}{2}\right)^2 \mathbb{I} \chi = \frac{\hbar^2}{4} \mathbb{I} \langle \chi^\dagger | \chi \rangle$
 $j = x, y, z$ $\xrightarrow{\text{forall } j} = \underline{\underline{\frac{\hbar^2}{4}}}$

$$\sigma_{S_x} = (\langle S_x^2 \rangle - \langle S_x \rangle^2)^{1/2} = \left(\frac{\hbar^2}{4} - 0\right)^{1/2} = \underline{\underline{\frac{\hbar}{2}}}$$

$$\sigma_{S_y} = (\langle S_y^2 \rangle - \langle S_y \rangle^2)^{1/2} = \left(\frac{\hbar^2}{4} - \frac{12^2 \hbar^2}{25^2}\right)^{1/2} = \hbar \left(\frac{624 - 576}{4 \times 25^2}\right)^{1/2} = \underline{\underline{\frac{7}{50} \hbar}}$$

$$\sigma_{S_z} = (\langle S_z^2 \rangle - \langle S_z \rangle^2)^{1/2} = \left(\frac{\hbar^2}{4} - \frac{7^2 \hbar^2}{4 \times 25^2}\right)^{1/2} = \hbar \left(\frac{25^2 - 49}{2500}\right)^{1/2} = \underline{\underline{\frac{12}{25} \hbar}}$$

(d) Uncertainty principle Check:

$$\sigma_{S_x} \sigma_{S_y} = \frac{7}{100} \hbar^2 \quad \text{vs} \quad \frac{\hbar^2}{2} |\langle S_z \rangle| = \frac{\hbar^2}{2} \left(\frac{7\hbar}{50}\right) \quad (\text{Equal! but OK})$$

$$\sigma_{S_y} \sigma_{S_z} = \frac{7}{50} \times \frac{12}{25} \hbar^2 \quad \text{vs} \quad \frac{\hbar^2}{2} |\langle S_x \rangle| = \frac{\hbar^2}{2} (0) \quad (\text{yes})$$

$$\sigma_{S_x} \sigma_{S_z} = \left(\frac{\hbar}{2}\right) \left(\frac{12\hbar}{25}\right) = \frac{12\hbar^2}{50} \quad \text{vs} \quad \frac{\hbar^2}{2} |\langle S_y \rangle| = \frac{\hbar^2}{2} \frac{7}{50} \quad (\text{yes})$$

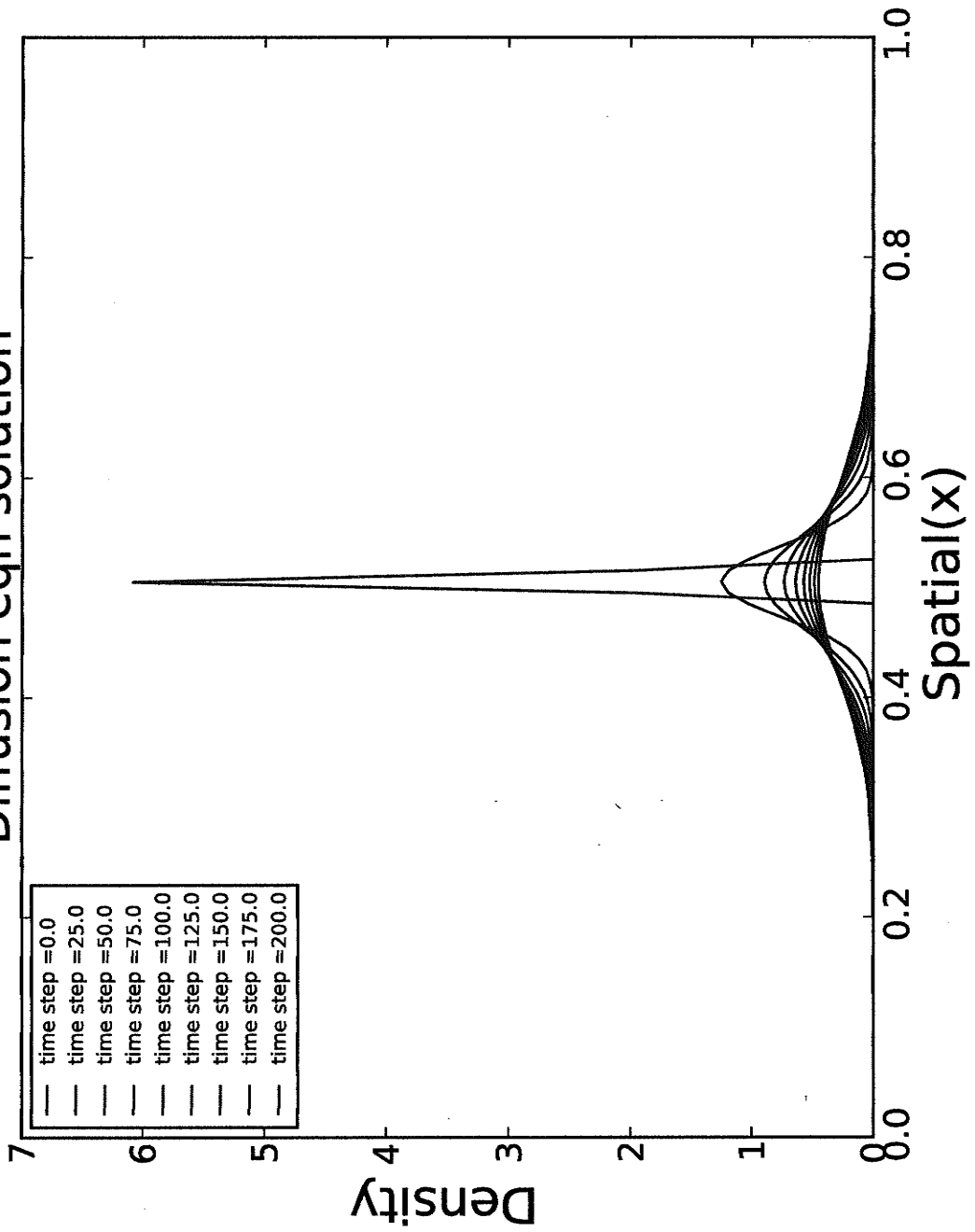
```
//
// P115B HW5 Prob 6
// g++ -o diff_eqn.exe diff_eqn.c
// Program to solve diffusion equation
// with forward time algorithm
#include <iostream>
#include <fstream>
#include <math.h>
using namespace std;
int main()
{
// Assign constants
//
const int N = 100; // Number of x-axis grid points
const int t_steps = 10; // Number of time steps
double L = 1.0; // X grid runs from 0 to L
double dx = L/(N-1); // Size of x-axis grid
double dt = 1.0e-6; // Size of time step
double D = 2.0; // D coeff in model
double D_factor = (D * dt)/(dx * dx);

// Initialize rho arrays
//
double rho[N], rhonew[N];
for (int i=0; i < N; i++)
{
rho[i] = 0.0;
rhonew[i] = 0.0;
}
rho[N/2] = 1.0/dx; // Initial condition delta function at center
// Set up disk files for output
ofstream fout("diffeqn_out1.dat");
//
// Loop over x grid and time steps
//
for (int i = 0; i < t_steps; i++) // Loop over time steps
{
for (int n = 1; n < N-1; n++) // Calculate new rho values
{
rhonew[n] = rho[n] + D_factor * ( rho[n+1] - 2*rho[n] + rho[n-1]);
}

for (int n = 0; n < N; n++) // Reset rho values for next time step & output
{
rho[n] = rhonew[n];
fout << i << " " << n << " " << rho[n]<< endl;
}
}

cout << " Finished writing data to file" << endl;
return 0;
}
```

Diffusion eqn solution



#6 Diffusion Equation Solver

$$D \frac{\partial^2 p}{\partial x^2} = \frac{\partial p}{\partial t} \Rightarrow p(x, t)$$

$$p(x+dx) = p(x) + p'(x) dx + \frac{1}{2} p''(x) dx^2$$

$$p(x-dx) = p(x) - p'(x) dx + \frac{1}{2} p''(x) dx^2$$

$$\Rightarrow p''(x) dx^2 = p(x+dx) + p(x-dx) - 2p(x)$$

$$\Rightarrow \frac{p(x, t+dt) - p(x, t)}{dt} = D \left[\frac{p(x+dx, t) + p(x-dx, t) - 2p(x, t)}{dx^2} \right]$$

$$= \left(\frac{\partial p}{\partial t} \right) = D \left(\frac{\partial^2 p}{\partial x^2} \right)$$

$$p(x, t+dt) = p(x, t) + \left(\frac{D \cdot dt}{dx^2} \right) [p(x+dx, t) + p(x-dx, t) - 2p(x, t)]$$

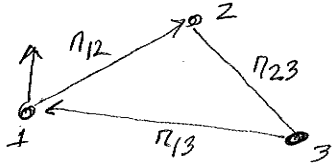
1D arrays are simpler / have lower memory requirement

2D arrays are easier if you want to create graphs/output in the same program (without writing to disk)

#7 Simulating 10^{11} stars in a galaxy

First, we're 3D so we have (3) position & (3) velocity components per star.

For 3 star system:



If position of star 1 moves from r_1 to r_1' , r_{12} & r_{13} must be recomputed.

In general, when looping through one time step for 10^{11} stars at star n , all stars from 1 to $n-1$ have moved which requires updates to $(n-1)$ pairwise π vectors and re-computes of acceleration & velocity changes for $(n-1)$ new vectors.

Let assume it takes 20 math ops to update 1 pairwise change. So to cycle through 10^2 stars would take

$$20 \text{ ops} \times \frac{n(n-1)}{2} \approx 10 \times 10^{2x} = 10^{2x+1} \text{ operation/cycle}$$

Now assume we want to update every year for 100 years (10^2 time cycles).

$$\text{Total compute time at } 10^9 \text{ ops/sec (flops)} = \frac{10^{2x+1} \text{ ops/cycle} \times 10^2 \text{ time cycle}}{10^9 \text{ ops/sec}} \approx 10^{2x-6} \text{ sec}$$

for $p=11 \Rightarrow 10^{22-6} = 10^{16}$ seconds $\approx 3 \times 10^8$ years to simulate 10^{11} stars

Let say, we only have 1-2 weeks of computer time in 10^6 sec.

$$\Rightarrow 2x-6 = 6 \Rightarrow x=6. \text{ We could model } 10^6 \text{ stars (million) } \textcircled{9}$$

With 10^5 processors (in parallel) $\frac{10^{2x+1} \times 10^2}{10^{14} \text{ ops/sec}} \approx 10^{2x-11} \approx 10^8 \text{ stars}$
 $2x \approx 17$ (100 million)