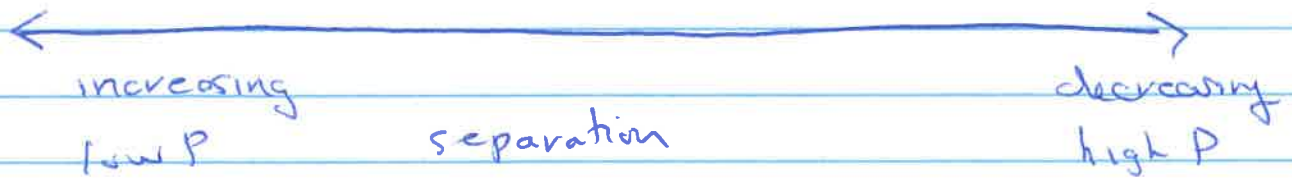


# Energy levels of $e^-$ in solid

## ① Physical picture

discrete atomic levels  
highly degenerate

broaden  
to bands



can even get metal-insulator transitions  
if bands themselves touch

## ② tight binding

← technically easy ←

## ③ solve schr. eqn

← technically hard



Bloch Thm  
underlies

must use  
computer  
"density  
functional  
theory"

but doesn't include  
details of nuclear  
potentials  
"chemistry"

Tight binding lattice of sites,  $e^-$  can hop between neighbors

$\psi(l)$  wave function on site  $l$

Amusing similarity to phonon math

lattice of atoms, springs between neighbors

$x(l)$  displacement of atom  $l$

$$H = \begin{pmatrix} E_a & -t & & \\ -t & E_a & -t & \\ & -t & E_a & -t \\ & & & \ddots \end{pmatrix} \quad \leftarrow d=1$$

$E_a = \text{original atomic level}$

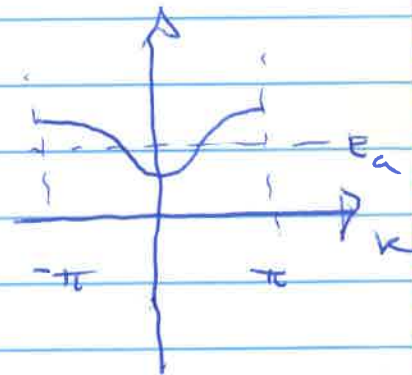
$$H\psi = E\psi$$

$$E_a \psi_l - t \psi_{l-1} - t \psi_{l+1} = E \psi_l$$

$$\psi_l = \psi_a e^{ikl}$$

$$E_a \psi_a + (-2t \cos k) \psi_a = E \psi_a$$

$$\text{If } \psi_a \neq 0 \quad E = E_a - 2t \cos k$$



Two types of atoms

$$\begin{array}{cccc}
 E_a & -t & & \\
 -t & E_b & -t & \\
 & -t & E_a & -t \\
 & & -t & E_b
 \end{array}$$

Method #1

$$\begin{array}{ll}
 \psi_l = \psi_a e^{ikl} & l = \text{odd} \\
 \psi_b e^{ikl} & l = \text{even}
 \end{array}$$

$$E_a \psi_a - 2t \cos k \psi_b = E \psi_a$$

$$E_b \psi_b - 2t \cos k \psi_a = E \psi_b$$

$$\begin{pmatrix} E_a - E & -2t \cos k \\ -2t \cos k & E_b - E \end{pmatrix} \begin{pmatrix} \psi_a \\ \psi_b \end{pmatrix} = 0$$



determinant must vanish

$$(E_a - E)(E_b - E) - (2t \cos k)^2 = 0$$

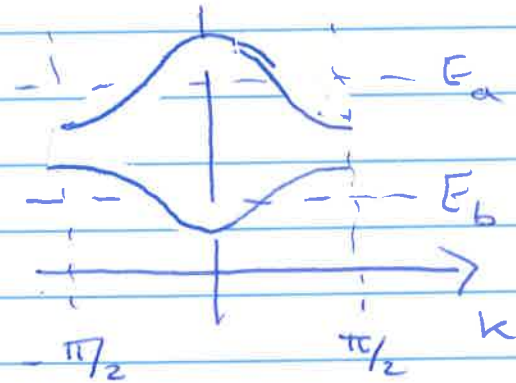
$$E^2 - (E_a + E_b)E + E_a E_b - (2t \cos k)^2 = 0$$

$$E = \frac{E_a + E_b \pm \sqrt{(E_a - E_b)^2 + (2t \cos k)^2}}{2}$$

$$t \text{ small} \quad (E_a - E_b) \left[ 1 + \left( \frac{2t \cos k}{E_a - E_b} \right)^2 \right]^{1/2}$$

$$\approx (E_a - E_b) \left[ 1 + \frac{1}{2} \left( \frac{2t \cos k}{E_a - E_b} \right)^2 \right]$$

$$E \approx \begin{cases} E_a + \frac{(t \cos k)^2}{E_a - E_b} \\ E_b - \frac{(t \cos k)^2}{E_a - E_b} \end{cases}$$



$$\text{NB} \quad \begin{aligned} \cos^2 k + \sin^2 k &= 1 \\ \cos^2 k - \sin^2 k &= \cos 2k \end{aligned}$$

$$\cos^2 k = \frac{1}{2} (1 + \cos 2k)$$

$$E \approx E_a + \frac{t^2}{2(E_a - E_b)} + \frac{t^2}{2(E_a - E_b)} \cos 2k$$

Already hint that  $E_a = E_b$  (degenerate)

is special. We will see more of this.



2D  
Square

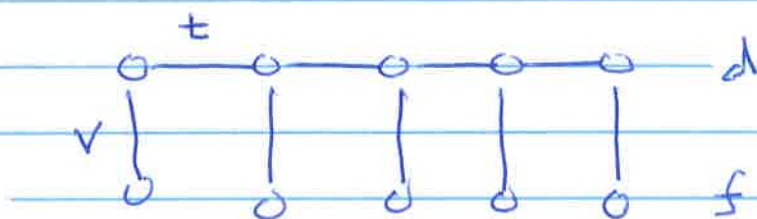
$$\psi(x, y) = \psi_0 e^{i(k_x x + k_y y)}$$

$$E_a \psi(x, y) - t \psi(x+1, y) - t \psi(x-1, y) - t \psi(x, y+1) - t \psi(x, y-1) = E \psi(x, y)$$

$$E = E_a - 2t \cos k_x - 2t \cos k_y$$

"PAM"

$d=1$   
for simplicity

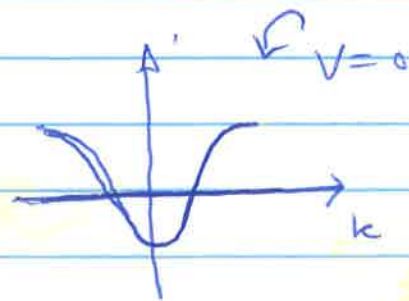


$$E_d \psi_d(x) - t \psi_d(x+1) - t \psi_d(x-1) - V \psi_f(x) = E \psi_d(x)$$

$$E_f \psi_f(x) - V \psi_d(x) = E \psi_f(x)$$

$$\begin{pmatrix} E_d - 2t \cos k & -V \\ -V & E_f \end{pmatrix} \begin{pmatrix} \psi_d \\ \psi_f \end{pmatrix} = E \begin{pmatrix} \psi_d \\ \psi_f \end{pmatrix}$$

Consider case  $E_d = E_f = 0$   
wlog



6.

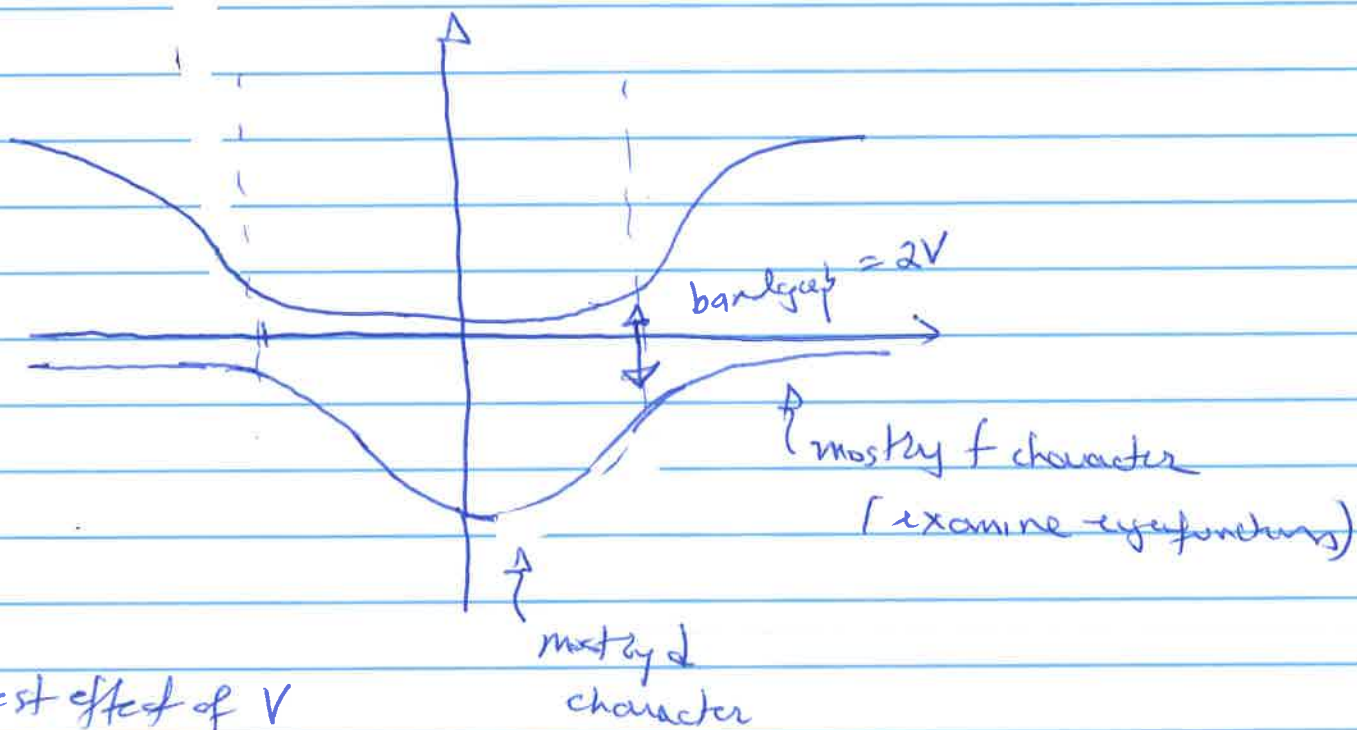
$$(-2t\cos k - E)(-E) - V^2 = 0$$

$$E^2 + E 2t\cos k - V^2 = 0$$

$$E = \left[ -2t\cos k \pm \sqrt{(2t\cos k)^2 + 4V^2} \right] / 2$$

If you sketch it:

$$k = \pi/2 \quad E = \pm V$$



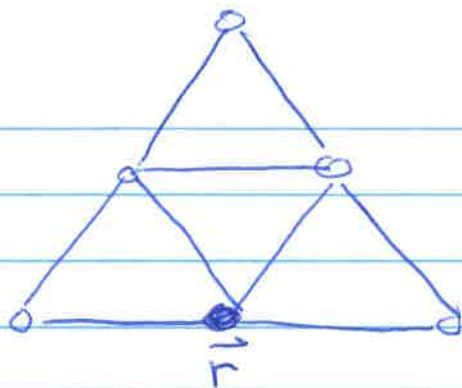
Biggest effect of  $V$   
is at degenerate  
points  $k = \pm \pi/2$   
where bands touch

mostly d  
character  
(eigenfunction)

Triangular lattice 2D

discrete indices less  
appropriate

square  $\vec{r} = l_x \hat{x} + l_y \hat{y}$



$$E_a \psi(\vec{r}) = -t \psi(\vec{r} + \hat{x}) - t \psi(\vec{r} + \frac{1}{2}\hat{x} + \frac{\sqrt{3}}{2}\hat{y})$$

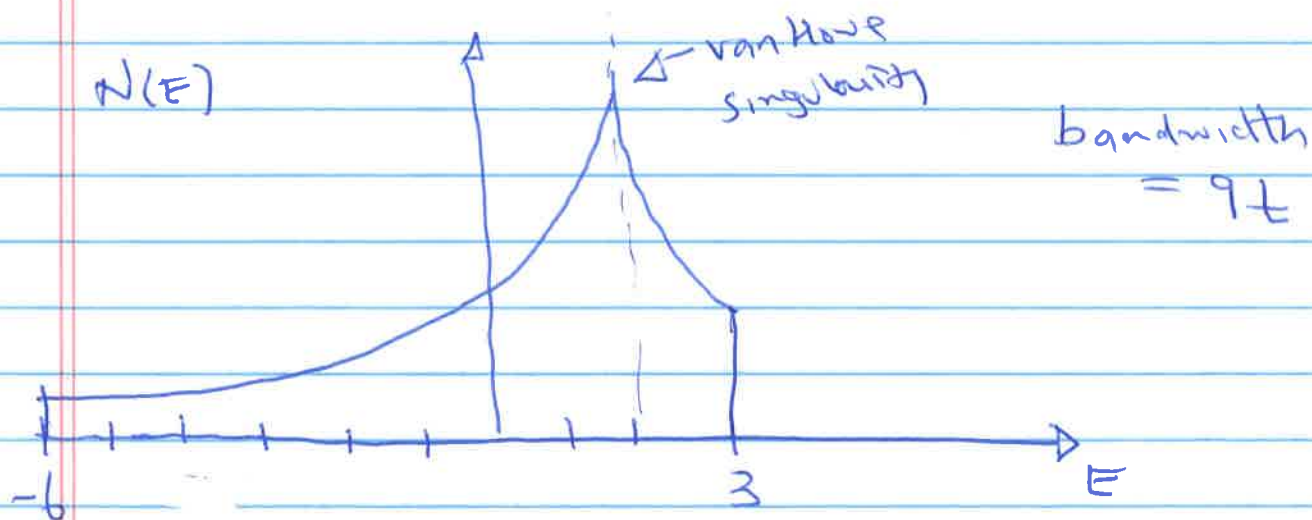
$$-t \psi(\vec{r} - \frac{1}{2}\hat{x} + \frac{\sqrt{3}}{2}\hat{y}) + hc$$

$$= E \psi(\vec{r})$$

$$\psi(\vec{r}) = \psi_a e^{i\vec{k} \cdot \vec{r}}$$

$\text{Na}_x\text{CoO}_2$   
Sodium cobaltate  
Co atoms on  
triangular lattice

$$E(k_x, k_y) = -2t \left[ \cos k_x + 2 \cos \frac{k_x}{2} \cos \frac{\sqrt{3}}{2} k_y \right]$$



"particle-hole symmetry"  $N(E) \neq N(-E)$