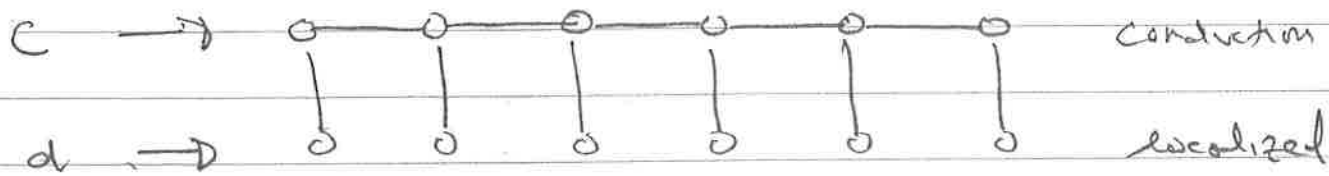


DAM-1

HW problem illustrated how hard it is to get ϵ_k for traditional approach. Let's see how easy it is in 2nd quantized



$$\hat{H} = -t \sum_l (c_{l+1}^\dagger c_l + c_l^\dagger c_{l+1}) + t' \sum_l (d_l^\dagger d_l + d_l^\dagger c_l)$$

$$c_l = \frac{1}{\sqrt{N}} \sum_k c_k e^{ikl}$$

$$d_l = \frac{1}{\sqrt{N}} \sum_k d_k e^{ikl}$$

$$\rightarrow \sum_k \epsilon_k c_k^\dagger c_k + t' \sum_l (\epsilon_k^\dagger d_l + d_l^\dagger c_k)$$

$$= \sum_k \begin{pmatrix} c_k^\dagger & d_k^\dagger \end{pmatrix} \begin{pmatrix} \epsilon_k & t' \\ t' & 0 \end{pmatrix} \begin{pmatrix} c_k \\ d_k \end{pmatrix}$$

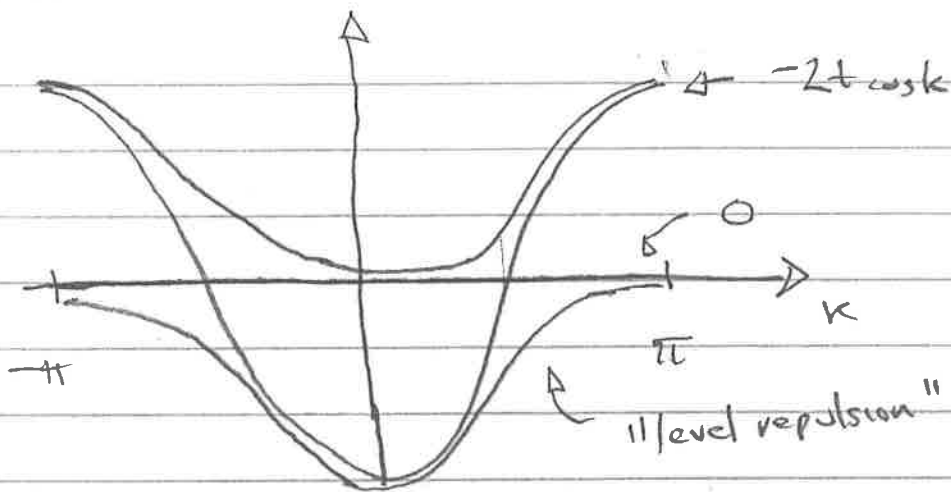
Diagonalize

$$(\epsilon_k - \lambda)(-\lambda) - t'^2 = 0$$

$$\lambda^2 - \epsilon_k \lambda - t'^2 = 0$$

$$\lambda = \frac{1}{2} \left[\epsilon_k \pm \sqrt{\epsilon_k^2 + 4t'^2} \right]$$

PAM-2



$$k=0 \quad \lambda = \frac{1}{2} \left[-2t \pm \sqrt{4t^2 + 4t'^2} \right] \approx -2t$$

$$+ 2t \left[1 + \frac{t'^2}{t^2} \right]^{1/2} \approx \frac{t'^2}{2t}$$

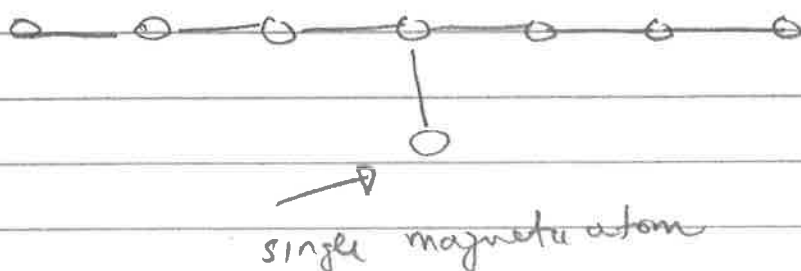
$$+ 2t \left[1 + \frac{t'^2}{2t^2} \right]$$

$$k = \pi/2 \quad \lambda = \frac{1}{2} \left[0 \pm \sqrt{4t'^2} \right] = \pm t'$$

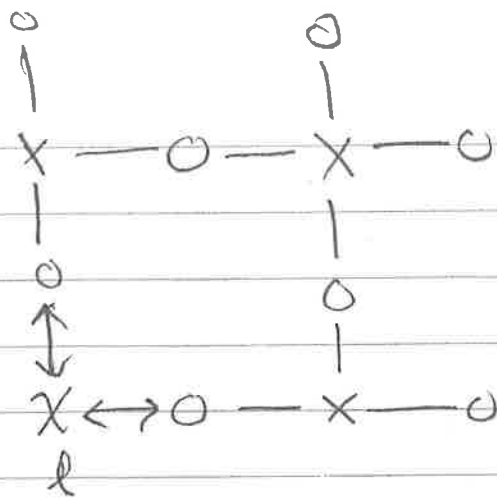
Note: We do not have any periodic potential

U_x applied here. Just 2 bands because 2 "orbitals" on each site

AIM as homework



$\text{CuO}_2 - 1$



$$\begin{aligned} H = & -t \sum_e \left[(d_e^\dagger P_{xl} + P_{xe}^\dagger d_e) \right. \\ & + (d_e^\dagger P_{yl} + P_{ye}^\dagger d_e) \\ & + (P_{xl}^\dagger d_{e+\hat{x}} + d_{e+\hat{x}}^\dagger P_{xe}) \\ & \left. + (P_{yl}^\dagger d_{e+\hat{y}} + d_{e+\hat{y}}^\dagger P_{ye}) \right] \end{aligned}$$

$$d_l = \frac{1}{\sqrt{N}} \sum_k e^{ikl} d_k$$

$$P_{xe} = \frac{1}{\sqrt{N}} \sum_k e^{ikl} P_{xk}$$

$$P_{yl} = \frac{1}{\sqrt{N}} \sum_k e^{ikl} P_{yk}$$

CuO₂-2

$$H = \frac{-t}{N} \sum_{\alpha} \sum_{\mathbf{k}} \sum_{\mathbf{k}'} \left\{ e^{-i\mathbf{k}\alpha} e^{i\mathbf{k}'\alpha} d_{\mathbf{k}}^{\dagger} p_{\mathbf{k}'}^{\dagger} + p_{\mathbf{k}}^{\dagger} d_{\mathbf{k}'} \right.$$

$$d_{\mathbf{k}}^{\dagger} p_{\mathbf{k}'}^{\dagger} + p_{\mathbf{k}}^{\dagger} d_{\mathbf{k}}$$

$$d_{\mathbf{k}}^{\dagger} e^{ik_x} p_{\mathbf{k}'}^{\dagger} + d_{\mathbf{k}} p_{\mathbf{k}'}^{\dagger} e^{-ik_x}$$

$$= \begin{pmatrix} d_{\mathbf{k}}^{\dagger} & p_{x\mathbf{k}}^{\dagger} & p_{y\mathbf{k}}^{\dagger} \end{pmatrix} \begin{bmatrix} 0 & -t(1+e^{ik_x}) & -t(1+e^{ik_y}) \\ -t() & 0 & 0 \\ -t() & 0 & 0 \end{bmatrix} \begin{pmatrix} d_{\mathbf{k}} \\ p_{x\mathbf{k}} \\ p_{y\mathbf{k}} \end{pmatrix}$$

\swarrow \searrow
 t_p

$$t\lambda(-\lambda)(-\lambda) + t(1+e^{ik_x})$$

~~$$-\lambda$$~~

$$-\lambda [(\epsilon_p - \lambda)(\epsilon_p - \lambda) - 0]$$

$$+ t(1+e^{ik_x}) [-t(1+e^{-ik_x})(t_p - \lambda)]$$

$$- t(1+e^{ik_y}) [+ t(1+e^{-ik_y})(t_p - \lambda)]$$

CuO₂-3

$$(\epsilon_p - \lambda) \left[-\lambda(\epsilon_p - \lambda) - t^2 (2 + 2\cos k_x + 2 + 2\cos k_y) \right]$$

$$\lambda = \epsilon_p$$

$$\lambda^2 - \lambda\epsilon_p - 2t^2 [1 + \cos k_x + \cos k_y] = 0$$

$$\lambda = \frac{1}{2} \left[\epsilon_p \pm \sqrt{\epsilon_p^2 + 8t^2 [1 + \cos k_x + \cos k_y]} \right]$$

If $\epsilon_p \gg t$



λ

$$\epsilon_p \sqrt{1 + \frac{8t^2}{\epsilon_p^2} (1 + \cos k_x + \cos k_y)}$$

$$\approx \epsilon_p \left[1 + \frac{4t^2}{\epsilon_p} (1 + \cos k_x + \cos k_y) \right]$$

$$\lambda_+ \approx \epsilon_p$$

$$\lambda_- \approx -\frac{2t^2}{\epsilon_p} - \frac{2t^2}{\epsilon_p} (\cos k_x + \cos k_y)$$



trivial
shift
in energy



2D square lattice ϵ_k !

$$t_{\text{eff}} \approx \frac{t^2}{\epsilon_p}$$

"2nd order
perturbation
theory"