

FSH-1

"Fine Structure" of Hydrogen

We saw spectrum of H atom

$$E_n = -\frac{me^4}{2\hbar^2} \frac{1}{n^2} = -\frac{e^2}{2a_0} \frac{1}{n^2} \quad a_0 = \frac{\hbar^2}{me^2}$$

And eigen functions

$$\Psi_{n\ell m}(r, \theta, \phi) = R_{n\ell}(r) Y_{\ell m}(\theta, \phi)$$

$$\Psi_{100} = \frac{1}{\sqrt{\pi a_0^3}} e^{-r/a_0}$$

Actually $m \rightarrow \mu$ reduced mass to allow for fact that nucleus is not fixed at origin.

"Fine structure" corrections $\propto (\alpha^2)$ $\alpha = \frac{e^2}{\hbar c} = \frac{1}{137}$

Two origins $\begin{cases} \rightarrow \text{relativistic correction} \\ \rightarrow \text{spin orbit coupling} \end{cases}$

FSH-2

We used $p^2/2m$ for kinetic energy, but really

$$\begin{aligned} & \sqrt{p^2 c^2 + m^2 c^4} - mc^2 \\ &= mc^2 \left\{ 1 + \frac{p^2}{m^2 c^2} \right\}^{1/2} - mc^2 \\ &= mc^2 \left\{ 1 + \frac{p^2}{2mc^2} - \frac{1}{8} \left(\frac{p^2}{m^2 c^2} \right)^2 + \dots \right\} - mc^2 \\ &= p^2/2m - p^4/8m^3 c^2 \end{aligned}$$

Use first order perturbation theory to evaluate

$$E_{n1} = -\frac{1}{8m^3 c^2} \langle \psi_{no} | \hat{p}^4 | \psi_{no} \rangle$$

A little trick is to write this as $\langle \psi_{no} | \hat{p}^2 \hat{p}^2 | \psi_{no} \rangle$

and use $\hat{p}^2 | \psi_{no} \rangle = 2m (E_{no} - V) | \psi_{no} \rangle$

from the Schroedinger Egn.

Then

$$\begin{aligned} E_{n1} &= -\frac{1}{2mc^2} \langle \psi_{no} | (E_{no} - V)^2 | \psi_{no} \rangle \\ &= -\frac{1}{2mc^2} \left[E_{no}^2 + 2E_{no} e^2 \langle \psi_{no} | \frac{1}{r} | \psi_{no} \rangle \right. \\ &\quad \left. + e^4 \langle \psi_{no} | \frac{1}{r^2} | \psi_{no} \rangle \right] \end{aligned}$$

FSH-3

These are some algebra to evaluate.

$$\langle \frac{1}{r} \rangle = \frac{1}{n^2 a} \quad \leftarrow \text{Easy}$$

$$\langle \frac{1}{r^2} \rangle = \frac{1}{(l + \frac{1}{2}) n^3 a^2} \quad \leftarrow \text{Hard!}$$

Derivation of $\langle \frac{1}{r} \rangle$:

FSH-4

Thus

$$E_{n1} = -\frac{1}{2mc^2} \left\{ E_{n0}^2 + 2E_{n0} \frac{e^2}{n^2 a} + \frac{e^4}{(l+1/2) n^3 a^2} \right\}$$

is the relativistic correction to H atom levels E_{n0} .

Note there is now an l dependence.

This can be rewritten as

$$E_{n1} = -\frac{E_{n0}^2}{2mc^2} \left[\frac{4n}{l+1/2} - 3 \right]$$

Q: What seems bogus about this calculation?!

A: We used non degenerate perturbation theory!

This was actually okay because the matrix elements of the particular perturbation between different degenerate H atom levels are zero, so matrix of V is diagonal.

FSM-5

Spin-orbit coupling

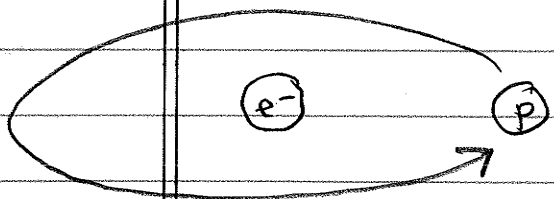
Q: Does anyone know what this is

A: e^- orbiting around nucleus is moving through \vec{E}

field of proton. From its point of view (its rest frame)

there is a \vec{B} field to which its spin couples

$$V = -\vec{\mu} \cdot \vec{B}$$



\vec{B} at center of current loop is a standard EM problem

Rest frame of e^-
has p circling
around

$$B = \frac{\mu_0 I}{2r} \quad \frac{2\pi I}{rc^2}$$

MKS CGS

$$I = \frac{e}{T}$$

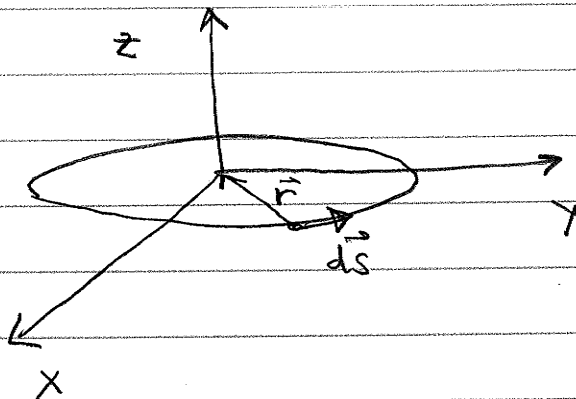
↑
period

$$L = mrv = 2\pi mr^2 / T$$

$$\text{so } I = e \sqrt{\frac{2\pi m^2 r^2}{L}} \quad \frac{\mu_0}{4\pi} \rightarrow 0.1$$

Biot-Savart

$$\vec{dB} = \frac{\mu_0}{4\pi} I \frac{d\vec{s} \times \vec{r}}{r^3}$$

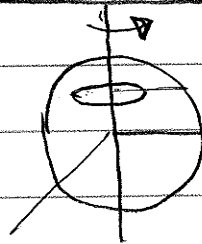


$$\vec{ds} \times \vec{r} = (r d\theta)(r) \hat{z}$$

$\uparrow \quad \uparrow$
 $|ds| \quad |r| \sin \frac{\pi}{2}$

$$\begin{aligned} \vec{B} &= \int d\vec{B} = \int_0^{2\pi} \frac{\mu_0}{4\pi} I \frac{r^2 d\theta}{r^3} \hat{z} \\ &= \frac{\mu_0 I}{2r} \hat{z} \end{aligned}$$

Spinning Ball of charge



$$\mu = IA$$

$$\mu = \int_{-r}^r dz \int_0^{\sqrt{r^2-z^2}} dp \cdot 2\pi p \cdot \frac{Q}{\frac{4}{3}\pi r^3} \cdot \underbrace{\pi p^2}_{A} \cdot \frac{1}{T}$$

$2\pi p dp dz$ is volume of ring
 charge density

$$\begin{aligned} &= \frac{3\pi Q}{2T r^3} \int_{-r}^r dz \int_0^{\sqrt{r^2-z^2}} p^3 dp = \frac{3\pi Q}{8T r^3} \int_{-r}^r dz (r^2-z^2)^2 \\ &= \frac{3\pi Q}{8T r^3} \left(r^4 z - \frac{2r^2 z^3}{3} + \frac{z^5}{5} \right) \Big|_{-r}^r = \frac{3\pi Q}{4T r^3} \left(1 - \frac{2}{3} + \frac{1}{5} \right) r^5 = \frac{2}{5} \frac{\pi r^2}{T} Q \end{aligned}$$

FSH-5''

Sphere moment of Inertia (basically same calculation)

$$I = \frac{2}{5} M r^2$$

$$L = I \omega = I \frac{2\pi}{T}$$

$$\mu = Q \frac{2}{5} \pi \frac{r^2}{T}$$

$$L = \frac{2}{5} M r^2 \frac{2\pi}{T}$$

We see that $\mu = \frac{2Q}{2M} L$

Actually the right answer is

$$\mu = -\frac{e}{M} S$$

This factor of 2 is Landé g-factor

and was only explained by Dirac!

FSH-6

$$\vec{B} = \frac{\mu_0}{2r} \frac{eL}{2\pi m r^2} = \frac{\mu_0}{4\pi} \frac{eL}{m r^3}$$

$$V = -\vec{\mu} \cdot \vec{B}$$

So need $\vec{\mu}$. If you do a ^{classical} spinning ball of charge you get

(see pages FSH-5', 5'')

$$\vec{\mu} = \frac{eQ}{2M} \vec{L}$$

Actual result is

"Lande g-factor"

$$\vec{\mu} = -\frac{e}{M} \vec{S}$$

Explained by Dirac Eqn

So ... $V_{SO} = -\vec{\mu} \cdot \vec{B}$

$$= \frac{e}{m} \vec{S} \cdot \frac{\mu_0}{4\pi} \frac{eL}{m r^3}$$

$$= \frac{\mu_0 e^2}{4\pi} \frac{1}{m^2 r^3} \vec{S} \cdot \vec{L}$$

One often writes this as

$$\vec{S} \cdot \vec{L} = \left[(\vec{L} + \vec{S})^2 - L^2 - S^2 \right] \frac{1}{2} = \frac{1}{2} [J^2 - L^2 - S^2]$$

\uparrow
 $\frac{1}{2}(\frac{1}{2} + 1)$

PS4-7

$$V_{so} = \frac{\mu_0 e^2}{8\pi} \frac{\hbar^2}{m^2 r^3} \left[j(j+1) - l(l+1) - 3/4 \right]$$

$$\left\langle \frac{1}{r^3} \right\rangle = \frac{1}{n^3 a^3} \frac{1}{l(l+1/2)(l+1)}$$

$$\left\langle \frac{1}{r} \right\rangle = \frac{1}{n^2} \frac{1}{a}$$

"Big Picture" Imagine you are a physicist in 1920's

QM is a new theory

You do these detailed calculations and they all
match up to expt!