

B-0

Bessel Functions (Laplacian in cylindrical coordinates)

$$\nabla^2 = \frac{\partial^2}{\partial z^2} + \frac{1}{\rho^2} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho}$$

Diffusion Eqn $\quad \mathcal{D} \nabla^2 \psi = \frac{\partial \psi}{\partial t}$

Try soln of form $\psi(\rho, \phi, z, t) = A(\rho) B(z) e^{\text{im}\phi} e^{-\alpha t}$

$\nearrow \quad \nearrow$
 We are experienced
 enough to guess the
 ϕ, t dependence...

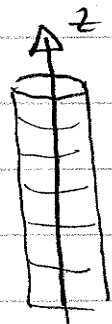
$$\mathcal{D} \left[\frac{\partial^2}{\partial z^2} + \frac{1}{\rho^2} \frac{\partial}{\partial \phi^2} + \frac{\partial^2}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial}{\partial \rho} \right] A B e^{\text{im}\phi} e^{-\alpha t} = -\alpha A B e^{\text{im}\phi} e^{-\alpha t}$$

$$\frac{1}{A} \left(\frac{\partial^2}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial}{\partial \rho} + \frac{(\alpha - m^2)}{\rho^2} \right) A(\rho) + \frac{1}{B} \frac{\partial^2}{\partial z^2} B(z)$$

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∞ long
cylinder
+ rotational
symmetry
as r
z axis

We will first consider case of no z, φ dependence

Since Bessel functions already appear there. Then

we will consider less symmetrical case

{ Schrodinger Egn
in "pillbox" ← class
Laplace Egn ← HW
in "pillbox"

$$D \nabla^2 \psi = \partial \psi / \partial t \quad \psi = A(\rho) e^{-\alpha t}$$

$$D \left(\frac{\partial^2}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial}{\partial \rho} \right) A e^{-\alpha t} = -\alpha A e^{-\alpha t}$$

$$\left(\frac{\partial^2}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial}{\partial \rho} + \frac{\alpha}{D} \right) A(\rho) = 0$$

$$\left[\rho^2 \frac{\partial^2}{\partial \rho^2} + \rho \frac{\partial}{\partial \rho} + \left(\frac{\alpha}{D} \right) \rho^2 \right] A(\rho) = 0$$

↑
"k²"

B-1

Bessel Functions

← soln of $\nabla^2 \psi = -k^2 \psi$
also stat mech!

Thus In solving diffusion Eqn for cylindrical geometry

we encounter

$$\left[x^2 \frac{d}{dx^2} + x \frac{d}{dx} + k^2 x^2 \right] u(x) = 0$$

These solns are a special case ($n=0$) of the Bessel Eqn

$$* \left[x^2 \frac{d^2}{dx^2} + x \frac{d}{dx} + (x^2 - n^2) \right] J_n(x) = 0$$

I.e. $u(x) = J_0(kx)$.

{ note that if $s \equiv kx$ $s \frac{d}{ds} = x \frac{d}{dx}$ and $s^2 \frac{d^2}{ds^2} = x^2 \frac{d^2}{dx^2}$

so that $(x^2 \frac{d^2}{dx^2} + x \frac{d}{dx} + k^2 x^2) u(kx) = s^2$

$$= (s^2 \frac{d^2}{ds^2} + s \frac{d}{ds} + s^2) u(s)$$

↳ Bessel for $n=0$ }

What are Bessel functions, the solns to * ?!

Basically they are like sines and cosines (oscillating)

except they also decay, and zeroes not evenly spaced

B-2

Sines and cosines can be defined in different ways.

Most simply $\frac{d^2}{dx^2} u(x) = -k^2 u$

But also

$\frac{d^2 u}{dx^2} = -k^2 u$
 $u = \sinh, \cosh$
 series?

$$\sin x = x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots$$

$$\cos x = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \frac{x^6}{6!} + \dots$$

• \sin, \cos complete \Rightarrow Fourier rep

$$\frac{1}{\pi} \int_0^{2\pi} \sin nx \sin mx dx = \delta_{nm}$$

$$\frac{1}{2\pi} \int_{-\omega}^{\omega} e^{i(k-k')x} dx = \delta(k-k')$$

For Bessel functions

$$J_n(x) = \sum_s \frac{1}{s! (s+n)!} (-1)^s \left(\frac{x}{2}\right)^{2s+n}$$

eg $J_0(x) = \sum_s \frac{1}{(s!)^2} (-1)^s \left(\frac{x}{2}\right)^{2s}$

$$= 1 - \frac{x^2}{4} + \frac{x^4}{4 \cdot 16} - \frac{x^6}{36 \cdot 64} + \dots$$

$$J_1(x) = \sum_s \frac{1}{s! (s+1)!} (-1)^s \left(\frac{x}{2}\right)^{2s+1}$$

$$= \frac{x}{2} - \frac{x^3}{16} + \frac{x^5}{96} - \dots$$

Where does this series expansion come from?

$$\frac{d}{dx} \sin x = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \dots$$

$$\frac{d^2}{dx^2} \sin x = -x + \frac{x^3}{3!} - \dots = -\sin x$$

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$$\frac{d}{dx} J_0(x) = -\frac{x}{2} + \frac{x^3}{16} - \frac{x^5}{6.64}$$

$$\frac{d^2}{dx^2} J_0(x) = -\frac{1}{2} + \frac{3}{16}x^2 - \frac{5x^4}{6.64}$$

$$x^2 \frac{d^3}{dx^3} J_0(x) + x \frac{d}{dx} J_0(x) + x^2 J_0(x)$$

$$= -\frac{1}{2}x^2 + \frac{3}{16}x^4 - \frac{5}{6.64}x^6$$

$$- \frac{x^2}{2} + \frac{x^4}{16} - \frac{x^6}{6.64}$$

$$+ x^2 - \frac{x^4}{4} + \frac{x^6}{4 \cdot 16} = 0$$

$$\begin{array}{ccc} \downarrow & \downarrow & \downarrow \\ 0 & 0 & 0 \end{array}$$

(1) Bessel DE

So

(2) SERIES

(3) other...

~~So it works but what does it come from~~

$$e^{x/2(t-1/t)} = \sum_n J_n(x) t^n$$

↑

generating function for Bessel

can show (1) gives series of page B-2

(2) satisfies Bessel DE of page B-1.

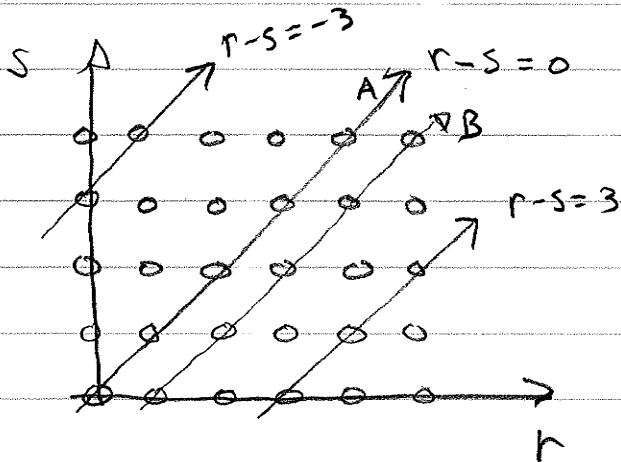
Basically a bunch of tedious algebra...

HW: Use $g(x, t)$ to prove $J_{n-1}(x) + J_{n+1}(x) = \frac{2n}{x} J_n(x)$

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Proof of (1)

$$e^{x/2} e^{-x/2} = \sum_{r=0}^{\infty} \left(\frac{x}{2}\right)^r \frac{t^r}{r!} \sum_{s=0}^{\infty} \left(\frac{x}{2}\right)^s (-1)^s \frac{t^{-s}}{s!}$$



$$t^{r-s} \quad n = r-s$$

$$s = r-n$$

$$r = s+n$$

can
relabel double sum
using r and s instead

n can take any value $-\infty$ to $+\infty$

If $n \geq 0$ s will start at \emptyset

$$\left(\sum_{n=-\infty}^{\infty} \sum_{s=-n}^{\infty} \right) + \sum_{n=0}^{\infty} \sum_{s=0}^{\infty} \left(\frac{x}{2}\right)^{s+n} \frac{t^{s+n}}{(s+n)!} \left(\frac{x}{2}\right)^s (-1)^s \frac{t^{-s}}{s!}$$

$n=0 \quad s=0,1,2,3,\dots \rightarrow$ line A since $r=s+n=s$
 $n=1 \quad s=0,1,2,3,\dots \rightarrow$ line B since $r=s+n=s+1$

Focus on
this to get
 J_n for $n > 0$

$$= \sum_{n=0}^{\infty} t^n \sum_{s=0}^{\infty} \frac{1}{s!} \frac{1}{(s+n)!} (-1)^s \left(\frac{x}{2}\right)^{2s+n}$$

Considering $n < 0$ in same way leads to

$$J_n(x) = -J_{-n}(x)$$

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$$g(x, t) = e^{x/2(t - 1/t)}$$

proof of (2)

$$\frac{d}{dx} g(x, t) = \frac{1}{2} \left(t - \frac{1}{t}\right) e^{x/2(t - 1/t)}$$

$$\frac{d^2}{dx^2} g(x, t) = \frac{1}{4} \left(t - \frac{1}{t}\right)^2 e^{x/2(t - 1/t)}$$

$$\therefore \left(x^2 \frac{d^2}{dx^2} + x \frac{d}{dx} + x^2 - n^2 \right) g(x, t)$$

$$= \left[x^2 \frac{1}{4} \left(t - \frac{1}{t}\right)^2 + x \frac{1}{2} \left(t - \frac{1}{t}\right) + x^2 - n^2 \right] e^{x/2(t - 1/t)}$$

$$= \left[x^2 \frac{1}{4} t^2 - \frac{x^2}{2} + \frac{x^2}{4t^2} + \frac{xt}{2} - \frac{x}{2t} + x^2 - n^2 \right] \quad "$$

111

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Further connections to families $\sin x$ and $\cos x$

$$\cos x = 1 - \frac{x^2}{2} + \frac{x^4}{24} - \frac{x^6}{720} \dots$$

$$= \sum_s \frac{1}{(2s)!} x^{2s} (-1)^s$$

$$\cos(ix) = \cosh x$$

$$\cosh x = 1 + \frac{x^2}{2} + \frac{x^4}{24} + \frac{x^6}{720} \dots$$

$$= \sum_s \frac{1}{(2s)!} x^{2s}$$

Similarly

$$J_n(x) = \sum \frac{1}{s!} \frac{1}{(s+n)!} \left(\frac{x}{2}\right)^{2s+n} (-1)^s$$

$$\rightarrow I_n(x) =$$

↓
eliminated

"Modified Bessel Function"

others Neumann, Hankel are essentially like

$$\cos x = \frac{e^{ix} + e^{-ix}}{2}$$

$$\sin x = \frac{e^{ix} - e^{-ix}}{2i}$$

$$\cos x, \sin x \quad \text{vs} \quad e^{ix}, e^{-ix}$$

B-7

An interesting identity comes from replacing t in generating function by $e^{i\theta}$

$$\frac{x}{2}(t + 1/t) \rightarrow x \cos \theta$$

$$e^{x \cos \theta} = \sum_{n=-\infty}^{\infty} J_n(x) e^{in\theta}$$

↑

We will see next quarter that this allows us to solve a very interesting statistical mechanics problem, "the XY chain" in terms of Bessel functions.