

Green's functions

We will begin in $d=2$ where we can draw nicer pictures + build intuition.

Solving a differential equation or partial differential equation involves not only specifying the eqn but also the boundary conditions.

$$m \frac{d^2 x}{dt^2} = -mg \quad (+) \quad x(t=0) \quad \dot{x}(t=0)$$

$$\Rightarrow x(t) = x_0 + v_0 t - \frac{1}{2} g t^2$$

$$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$$

$$f(x, t) = A(x) B(t)$$

$$A \frac{dB}{dt} = D B \frac{d^2 A}{dx^2} \quad \frac{1}{D} \frac{1}{B} \frac{dB}{dt} = \frac{1}{A} \frac{d^2 A}{dx^2} = -k^2$$

The sign was incorporating some intuition about boundary conditions

$$\frac{dB}{dt} = -D k^2 B \quad B(t) = e^{-D k^2 t}$$

We could have chosen $+k^2$ in principle, but sol'n would grow exponentially in time.

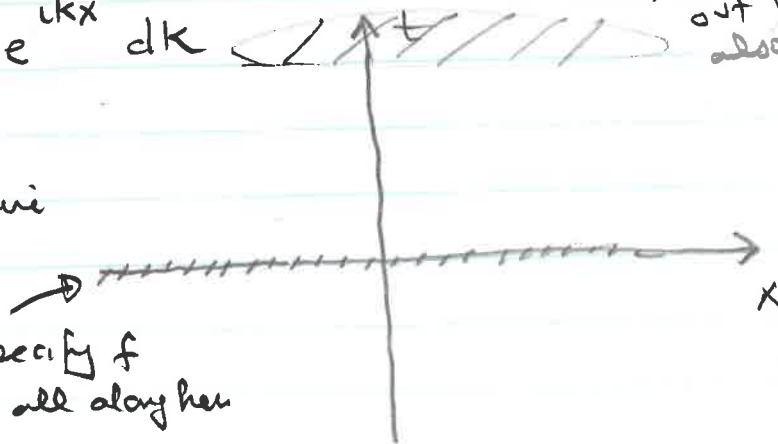
Our sol'n also involved a knowledge of $f(x, 0)$

$$f(x, t) = \int c(k) e^{-D k^2 t} e^{ikx} dk$$

need $f(x, 0)$ to determine

specify f all along here

specify f out here also



(2)

(2)

Consider wave eqn for vibrating string



$$\frac{\partial^2 y}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 y}{\partial t^2} = 0$$

$$y = A(x) B(t)$$

$$B \frac{d^2 A}{dx^2} = \frac{1}{v^2} A \frac{d^2 B}{dt^2}$$

$$\frac{1}{A} \frac{d^2 A}{dx^2} = \frac{1}{v^2} \frac{1}{B} \frac{d^2 B}{dt^2} = -k^2$$

again physics

$$A = \sin kx \quad \cos kx \quad B = \sin kv t \quad \cos kv t$$

more boundary conditions: $A(0) = 0 \quad B(L) = 0$

discard $\cos kx$ solns

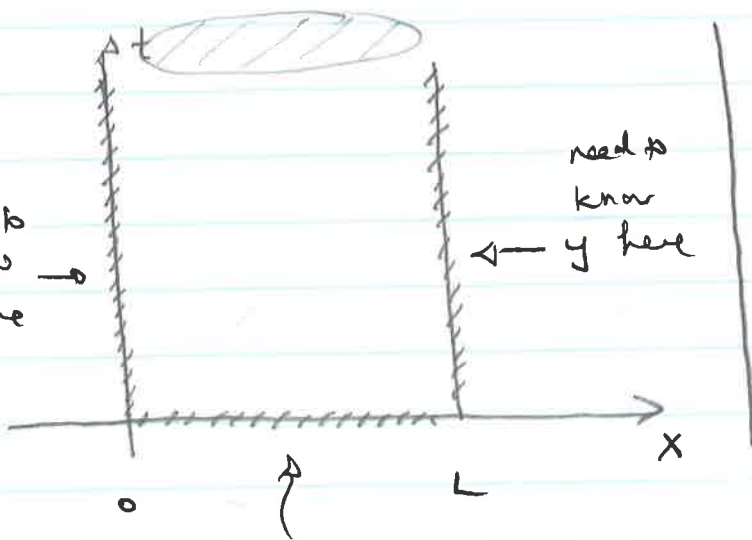
and restrict $k \Rightarrow k_n = \frac{\pi}{L} n$

$$y(x,t) = \sum_n \sin \frac{\pi x}{L} n \left(a_n \cos \frac{\pi t n v}{L} + b_n \sin \frac{\pi t n v}{L} \right)$$

Wave eqn needs open "u" shaped bdy

need to know y here

need to know y here



need to know y and $\frac{dy}{dt}$ here

$$y(x,0) = \sum_n a_n \sin \frac{n\pi x}{L} \quad a_n = \frac{2}{L} \int_0^L \sin \frac{n\pi x}{L} y(x,0) dx'$$

$$y(x,t) = \sum_n \frac{2}{L} \int_0^L \sin \frac{n\pi x'}{L} y(x',0) dx' \sin \frac{n\pi x}{L} \cos \frac{n\pi t v}{L}$$

$$y(x,0) = \int_0^L \sin \dots$$

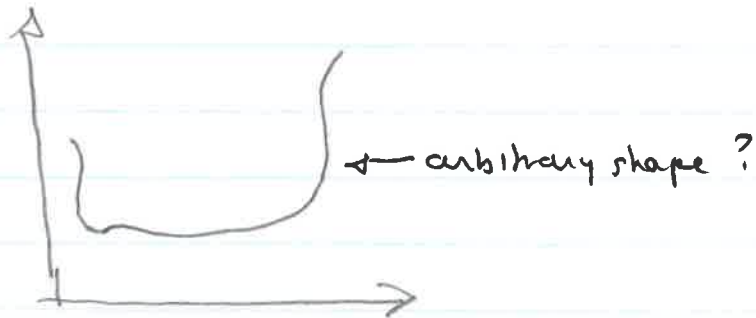
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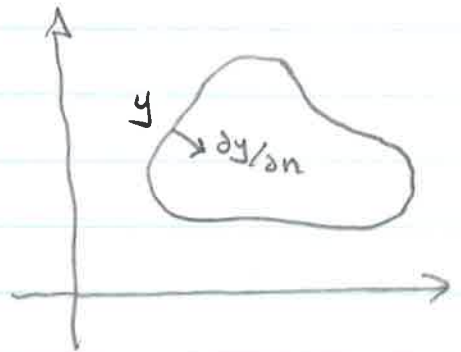
"Cauchy bdy conditions" : Specify value of function and its normal derivative ($\frac{\partial y}{\partial n}$ is change of y \perp to boundary at $t=0$)

Questions:

(1) Can you use more bizarre boundaries



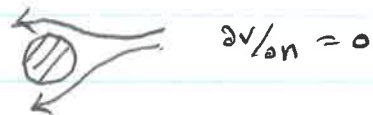
(2) What if curve is closed
would specifying y and $\frac{\partial y}{\partial n}$
everywhere be overly restrictive?



Ans: yes! only need y or $\frac{\partial y}{\partial n}$
if bdy closed

y : Dirichlet \leftarrow Laplace eqn with y specified

$\frac{\partial y}{\partial n}$: Neumann \leftarrow fluid flow around solid bodies

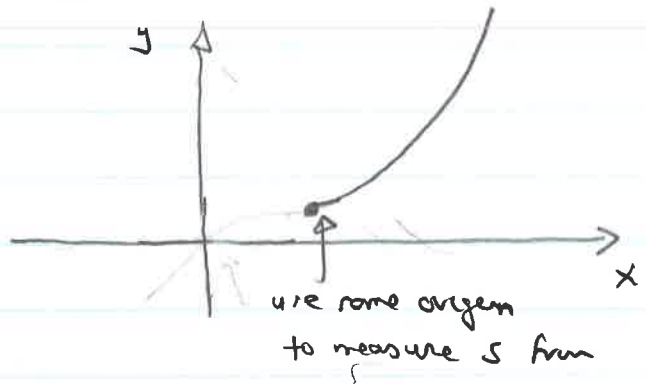


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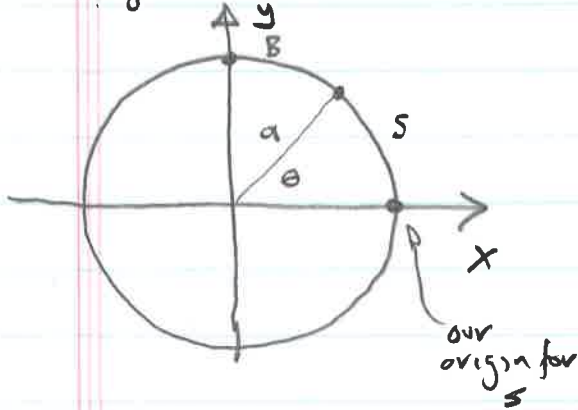
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How to express a general boundary?

$$x = \xi(s) \quad y = \eta(s)$$



Eg



$$x = a \cos \theta = a \cos\left(\frac{s}{a}\right) \equiv \xi(s)$$

$$y = a \sin \theta = a \sin\left(\frac{s}{a}\right) \equiv \eta(s)$$

Eg at "B" $s = \frac{2\pi}{4}a = \frac{\pi}{2}a, \theta = \frac{\pi}{2}$

Let's try computing $\psi(x, y)$ some distance from the body at (ξ, η)

$$\psi(x, y) = \psi(\xi, \eta) + (x - \xi) \frac{\partial \psi}{\partial x} + (y - \eta) \frac{\partial \psi}{\partial y}$$

$$+ \frac{1}{2} \left[(x - \xi)^2 \frac{\partial^2 \psi}{\partial x^2} + 2(x - \xi)(y - \eta) \frac{\partial^2 \psi}{\partial x \partial y} + (y - \eta)^2 \frac{\partial^2 \psi}{\partial y^2} \right]$$

If we knew all the partial derivatives then we would know ψ

We shall not need all of them, as our examples have

already hinted.

what do we know?

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simple analogy

Is this analogous?

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$$m \frac{d^2 x}{dt^2} = -mg$$

$$x(t) = x(0) + \left. \frac{dx}{dt} \right|_0 t + \frac{1}{2} \left. \frac{d^2 x}{dt^2} \right|_0 t^2 + \frac{1}{6} \left. \frac{d^3 x}{dt^3} \right|_0 t^3 + \dots$$

Need to specify $x|_0, \left. \frac{dx}{dt} \right|_0, \left. \frac{d^2 x}{dt^2} \right|_0, \dots$!?

Why do we not need higher ones?

The partial differential eqn gives the $\left. \frac{d^2 x}{dt^2} \right|_{t=0} = -g$

all higher ones zero!

Something similar here

If we are told ψ along the curve $[\xi(s), \eta(s)]$ then we can always compute $\psi(s) = \psi[\xi(s), \eta(s)]$

$$\frac{d\psi}{ds} = \frac{\partial \psi}{\partial x} \frac{d\xi}{ds} + \frac{\partial \psi}{\partial y} \frac{d\eta}{ds} \quad \leftarrow \text{one eqn relating } \frac{\partial \psi}{\partial x} \text{ and } \frac{\partial \psi}{\partial y}$$

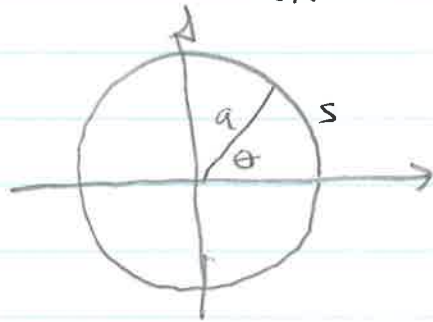
Do later @ *

Example! $\psi = 0$ on circle

$$\frac{d\xi}{ds} = -\sin s/a$$

$$d\eta/ds = +\cos s/a$$

$$-\frac{\partial \psi}{\partial x} \sin s/a + \frac{\partial \psi}{\partial y} \cos s/a = 0$$

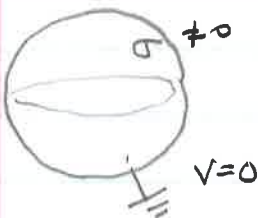


~~one eqn relating~~

suppose also we are told normal derivative of ψ is a nonzero constant σ

3-d analogy

$E \sim \sigma$
 $\nabla \cdot \vec{V}$

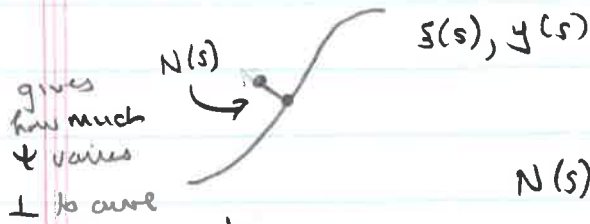


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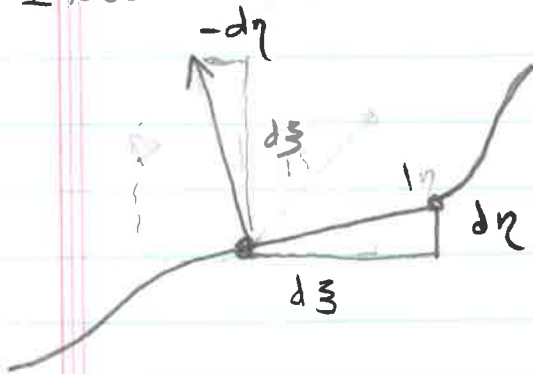
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If we are also told how ψ changes \perp to curve

Denote $N(s)$ as rate of change of ψ in direction \perp to curve



$$N(s) = \left(\frac{\partial \psi}{\partial x} \frac{dx}{ds} - \frac{\partial \psi}{\partial y} \frac{dy}{ds} \right)$$



normal is $-d\eta$

solve these 2 eqns for $\frac{\partial \psi}{\partial x}$ and $\frac{\partial \psi}{\partial y}$!

$$\begin{aligned} * \quad & + \frac{\partial \psi}{\partial x} \cos s/a + \frac{\partial \psi}{\partial y} \sin s/a = \sigma \quad \cos \quad \sin \\ & - \frac{\partial \psi}{\partial x} \sin s/a + \frac{\partial \psi}{\partial y} \cos s/a = 0 \quad -\sin \quad +\cos \end{aligned}$$

$$\frac{\partial \psi}{\partial x} = \sigma \cos s/a$$

$$\frac{\partial \psi}{\partial y} = \sigma \sin s/a$$

$$\frac{d\psi}{ds} = \frac{\partial \psi}{\partial x} \frac{dx}{ds} + \frac{\partial \psi}{\partial y} \frac{dy}{ds} \quad \frac{dx}{ds} \quad \frac{dy}{ds}$$

$$N(s) = \frac{\partial \psi}{\partial x} \frac{d\eta}{ds} - \frac{\partial \psi}{\partial y} \frac{dx}{ds} \quad + \frac{d\eta}{ds} \quad - \frac{dx}{ds}$$

$$\frac{\partial \psi}{\partial x} = N(s) \frac{d\eta}{ds} + \frac{dx}{ds} \frac{d\psi}{ds} \equiv p(s)$$

$$\frac{\partial \psi}{\partial y} = -N(s) \frac{dx}{ds} + \frac{d\eta}{ds} \frac{d\psi}{ds} \equiv q(s)$$

$$\text{NB } \left(\frac{dx}{ds} \right)^2 + \left(\frac{d\eta}{ds} \right)^2 = 1 !$$

known functions since
 $N(s), \frac{\partial \psi}{\partial x}, \eta, \xi$ all known

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Getting 2nd derivatives is crucial. If we get 2nd derivs, higher ones just "move of the same"

Three "unknowns" $\frac{\partial^2 \psi}{\partial x^2}$; $\frac{\partial^2 \psi}{\partial x \partial y}$; $\frac{\partial^2 \psi}{\partial y^2}$

2 Eqns immediately

$$x = z(s)$$

Remember $p(s)$ is known!

$$\frac{dp}{ds} = \frac{d}{ds} \frac{\partial \psi}{\partial x} = \frac{\partial^2 \psi}{\partial x^2} \frac{dz}{ds} + \frac{\partial^2 \psi}{\partial x \partial y} \frac{dy}{ds}$$

$$\frac{dq}{ds} = \frac{\partial^2 \psi}{\partial x \partial y} \frac{dz}{ds} + \frac{\partial^2 \psi}{\partial y^2} \frac{dy}{ds}$$

X

What's 3rd eqn?

What haven't we used at all yet?

The PDE

$$A(x,y) \frac{\partial^2 \psi}{\partial x^2} + 2B(x,y) \frac{\partial^2 \psi}{\partial x \partial y} + C(x,y) \frac{\partial^2 \psi}{\partial y^2} = F$$

$$F(x,y, \psi, \frac{\partial \psi}{\partial x}, \frac{\partial \psi}{\partial y})$$

If Eqn is linear $H(x,y) + D(x,y) \frac{\partial \psi}{\partial x}$

$+ E(x,y) \frac{\partial \psi}{\partial y} + G(x,y) \psi$

$$\begin{pmatrix} \frac{dz}{ds} & \frac{dy}{ds} & 0 \\ 0 & \frac{dz}{ds} & \frac{dy}{ds} \\ A & 2B & C \end{pmatrix} \begin{pmatrix} \frac{\partial^2 \psi}{\partial x^2} \\ \frac{\partial^2 \psi}{\partial x \partial y} \\ \frac{\partial^2 \psi}{\partial y^2} \end{pmatrix} = \begin{pmatrix} dp/ds \\ dq/ds \\ F(s) \end{pmatrix}$$

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Can show all higher derivatives can be computed in terms of known quantities.

* One caveat: Determinant $\neq 0$!

$$\begin{aligned} \frac{d^2 \xi}{ds^2} \left[c \frac{d\xi}{ds} - 2B \frac{d\eta}{ds} \right] - \frac{d\xi}{ds} \left[-A \frac{d\eta}{ds} \right] \\ = c \left(\frac{d\xi}{ds} \right)^2 - 2B \frac{d\xi}{ds} \frac{d\eta}{ds} + A \left(\frac{d\eta}{ds} \right)^2 \end{aligned}$$

$$d\xi = dx \quad d\eta = dy$$

$$c(x,y) dx^2 - 2B(x,y) dx dy + A(x,y) (dy)^2 = 0$$

Cauchy uniquely specifies soln unless the curve

$\xi(s), \eta(s)$ is a "characteristic" \rightarrow cont'd

It is interesting to understand why ~~to~~ giving ψ on N along a characteristic is insufficient.

One can discuss this problem generally, but maybe it's better to take the example of the wave eqn

$$\frac{\partial^2 \psi}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2} = 0$$

$$\begin{aligned} A(x,t) &= 1 \\ B(x,t) &= 0 \\ C(x,t) &= -1/c^2 \end{aligned}$$

$$-1/c^2 dx^2 + dt^2 = 0$$

$$\left(\frac{dx}{dt} \right)^2 = c^2$$

$$\frac{dx}{dt} = \pm c$$

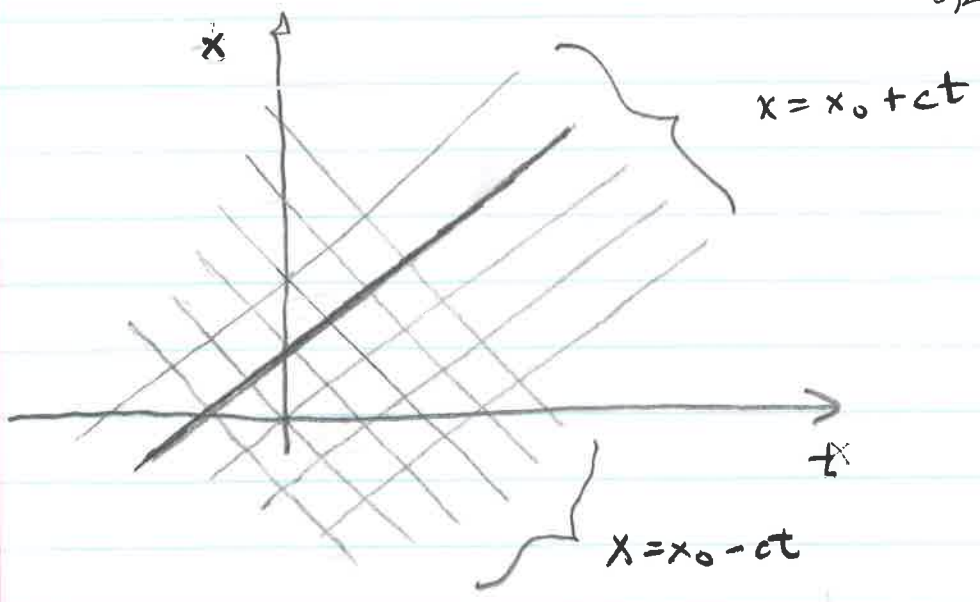
$$x = x_0 \pm ct$$

The solution to the wave eqn is the sum of

$$\psi(x,t) = f(x+ct) + g(x-ct)$$

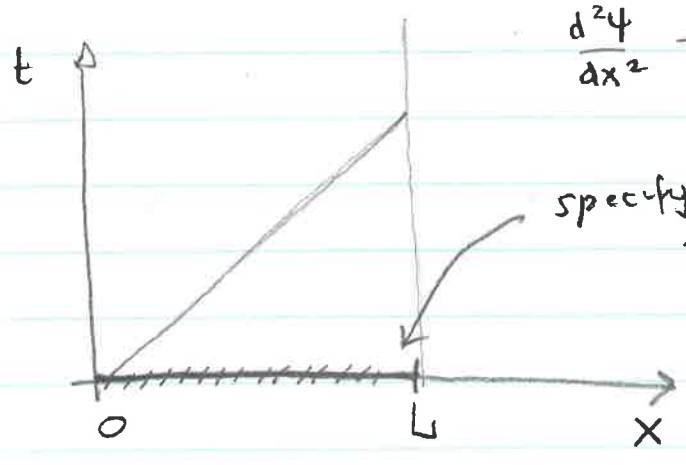
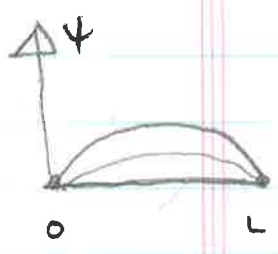
↑ ↑
arbitrary functions

Giving ψ along $x = x_0 \pm ct$ and $N(S)$ along there will only specify one of f or g and not the other



one needs a bdy which cuts across

More complex pdes may have/will have more complicated characteristics.



$$\frac{d^2\psi}{dx^2} - \frac{1}{c^2} \frac{d^2\psi}{dt^2} = 0$$

specifying ψ and $\frac{d\psi}{dt}$ here is good

Actually, there are always two families of characteristics:

$$A dy^2 - 2B dx dy + c dx^2 = 0$$

$$dy = \left[2B dx \pm \sqrt{4B^2 dx^2 - 4Ac dx^2} \right] / 2A$$

$$A dy = \left[B dx \pm \sqrt{B^2 - Ac} \right] dx$$

$$\frac{dy}{dx} = \frac{B \pm \sqrt{B^2 - Ac}}{A}$$

wave eqn $A = 1$ $C = -1/c^2$ $B = 0$

$$\frac{dt}{dx} = \frac{0 \pm \sqrt{1/c^2}}{1} = \pm 1/c$$

(10)

very explicitly:

(11)

$$\psi(x,t) = \sum_n \sin \frac{n\pi x}{L} \left(a_n \cos \frac{n\pi vt}{L} + b_n \sin \frac{n\pi vt}{L} \right)$$

$$= \sum_n \frac{a_n}{2} \left\{ \sin \frac{n\pi}{L} (x+vt) - \sin \frac{n\pi}{L} (x-vt) \right\}$$

$$2 \sin A \cos B$$

$$= \sin(A+B) - \sin(A-B)$$

$$+ \frac{b_n}{2} \left\{ \cos \frac{n\pi}{L} (x-vt) - \cos \frac{n\pi}{L} (x+vt) \right\}$$

$$2 \sin A \sin B$$

$$= \cos(A-B) - \cos(A+B)$$

$$= \sum_n \left[\frac{a_n}{2} \sin \frac{n\pi}{L} (x+vt) - \frac{b_n}{2} \cos \frac{n\pi}{L} (x+vt) \right] \leftarrow f(x+vt)$$

$$+ \left[-\frac{a_n}{2} \sin \frac{n\pi}{L} (x-vt) + \frac{b_n}{2} \cos \frac{n\pi}{L} (x-vt) \right] \leftarrow g(x-vt)$$

given $\psi(x,t)$ along $x = x_0 + vt$ only $x - vt = x_0 \dots$

(?)

~~Rotation type~~

characteristics also taught us that $B^2 - AC$ is an imp't quantity

$B(x,y)^2 > A(x,y)C(x,y)$ everywhere : hyperbolic

$0 > B^2 - AC$ (-1/c^2) wave eq'n

$AC > B^2$ elliptic

$AC = B^2$ (everywhere) parabolic

Diffusion

$\frac{\partial \psi}{\partial t} = D \frac{\partial^2 \psi}{\partial x^2}$

Laplace

$\frac{d^2 \psi}{dx^2} + \frac{d^2 \psi}{dy^2}$

$A=1 \quad B=1 \quad C=0$

Bottom Line

Conditions	Bdy	Hyperbolic	Elliptic	Parabolic
Dirichlet or Neumann (value or slope)	OPEN	insufficient	insufficient	unique sufficient
	CLOSED	sufficient	sufficient	overspecified
	OPEN	sufficient	sufficient	overspecified
Cauchy value and slope	CLOSED	overspecified	overspecified	overspecified

Now finally Green's functions...

We will consider $d=2$ Laplace Egn

$$\frac{d^2\psi}{dx^2} + \frac{d^2\psi}{dy^2} = 0$$

$$\psi = \bar{X}(x) \bar{Y}(y)$$

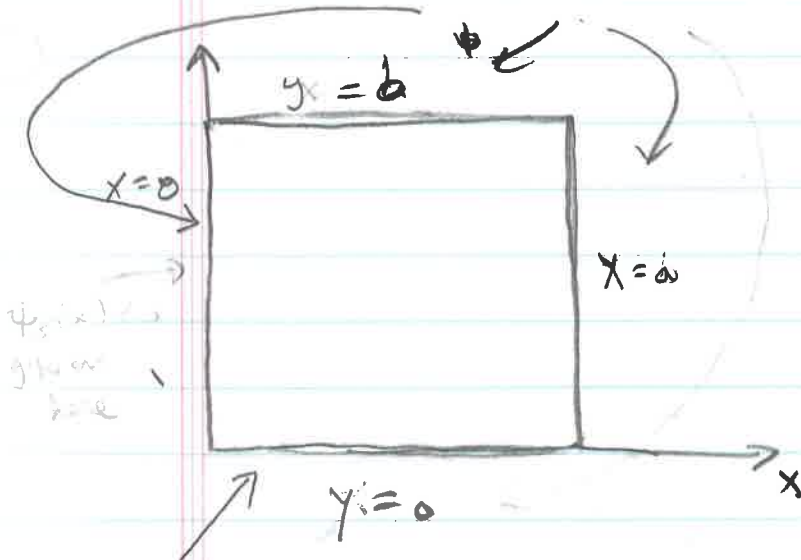
$$\frac{\bar{X}''}{\bar{X}} = -k^2 = -\frac{\bar{Y}''}{\bar{Y}}$$

$\sin kx$
 $\cos kx$

$\sinh ky$
 $\cosh ky$

Elliptic: we claimed
Dirichlet or
Neumann on
closed surface
sufficient

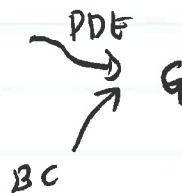
Consider $\psi(x,y) = 0$ along $x=0$ $x=a$ $y=b$



$\psi(x,0) \neq 0$
here

We will construct the Green function
for this PDE and these BC.

Emphasise



$x=0, x=a$

can guarantee b.c by choosing $k = \frac{n\pi}{a}$ and

taking only the sine solutions

$$\psi(x,y) = \sum b_n \sin \frac{n\pi}{a} x \sinh \frac{n\pi}{a} (b-y)$$

Also can guarantee $xy=b$ b.c by choosing $\sinh(b-y)k$

$$\begin{aligned} \left[\sinh(A+B) &= \frac{1}{2} [e^{A+B} - e^{-A-B}] \right. \\ &= \frac{1}{4} (e^A - e^{-A})(e^B + e^{-B}) + \frac{1}{4} (e^A + e^{-A})(e^B - e^{-B}) \\ &= \sinh A \cosh B + \cosh A \sinh B \left. \right] \end{aligned}$$

$$\psi_s(x) \equiv \psi(x,y=0) = \sum b_n \sin \frac{n\pi x}{a} \sinh \frac{n\pi b}{a}$$

$$\int_0^a dx \sin \frac{n\pi x}{a} \psi_s(x) = \frac{a}{2} b_n \sinh \frac{n\pi b}{a}$$

$$\begin{aligned} \int_0^a \sin^2 \frac{n\pi x}{a} dx \\ \Rightarrow \theta = \frac{n\pi x}{a} \\ \int_0^{n\pi} \sin^2 \theta d\theta \frac{a}{n\pi} \\ = \frac{n\pi}{2} \frac{a}{n\pi} = \frac{a}{2} \end{aligned}$$

$$\psi(x,y) = \sum_{n=1}^{\infty} \left(\frac{2}{a} \int_0^a dx' \sin \frac{n\pi x'}{a} \psi_s(x') \right) \sin \frac{n\pi x}{a} \frac{\sinh \frac{n\pi}{a} (b-y)}{\sinh \frac{n\pi b}{a}}$$

~~Consider limits $b \rightarrow \infty$ and $a \rightarrow \infty$ (solve Laplace Eqn in whp given ψ along $x=0$)~~

$$\begin{aligned} \psi(x,y) &= \int dx' \psi_s(x') \underbrace{f(x-x',y)} \\ &= \frac{2}{a} \sum_{n=1}^{\infty} \frac{\sinh \frac{n\pi}{a} (b-y)}{\sinh \frac{n\pi b}{a}} \left(\sin \frac{n\pi x}{a} \sin \frac{n\pi x'}{a} \right) \end{aligned}$$