

PHYSICS 110A, WINTER 2017
ELECTRICITY AND MAGNETISM

Assignment Seven, Due Friday, March 9, 5:00 pm.

[1.] Griffiths 3-24.

[2.] An infinite cylinder of radius $r = a$ has charge density $\sigma(\phi) = \sigma_0 \cos 4\phi$ on its surface. Compute the potential $V(\rho, \phi, z)$ inside (i.e. for $\rho < a$) and also outside ($\rho > a$).

[3.] Extra credit: Go through separation of variables solution of Laplace's equation in cylindrical coordinates when it is *not* legitimate to ignore the z dependence. (This is an extension of Problem [1].) Determine what the functional form is for the z and ϕ dependence. Look up online what the solutions are for the ρ dependence, just so you know in the future what those functions are.

[4.] Compute the monopole, dipole, and quadrupole terms of the potential of a point charge q located at position $\mathbf{r} = (0, 0, a)$. What is true for the case $a = 0$?

[5.] Compute the monopole, dipole, and quadrupole terms of the potential of two point charges, $+q$ located at position $\mathbf{r} = (0, 0, a)$, and $-q$ located at position $\mathbf{r} = (0, 0, b)$. Comment on the dependence of the dipole term on a, b . In particular, if you shift the charges by the same amount in the z direction, to $(0, 0, a + c)$ and $(0, 0, b + c)$, what happens to the dipole term? Enunciate a general rule about the monopole and dipole terms, the nature of the charge distribution, and the choice of origin.

Physics 110A Winter 2018
Assignment 7

1

See the solution to problem #2, where

we conclude solns are of the form

$$\rho^n \cos n\phi; \rho^n \sin n\phi; \rho^{-n} \cos n\phi, \rho^{-n} \sin n\phi$$

where $n=1, 2, 3, \dots$ is a positive integer.

To this list we must also add the case where $n=0$ so there is no ϕ dependence and

$$\rho \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} R(\rho) = 0$$

This has the soln $R(\rho) = C$, a constant, but also $R(\rho) = \ln \rho$ since

$$\rho \frac{\partial}{\partial \rho} \ln \rho = \rho \frac{1}{\rho} = 1$$

so that when we take the second $\rho \frac{\partial}{\partial \rho}$ in the Laplacian we get zero. This $\ln \rho$ term is of course the potential due to an infinite line of charge.

2 In cylindrical coordinates (ρ, ϕ, z)

$$\nabla^2 V = \frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial V}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2}$$

If there is no z dependence, and we write $V(\rho, \phi) = R(\rho) Q(\phi)$ then Laplace's eqn $\nabla^2 V = 0$ becomes

$$Q \frac{1}{\rho} \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} R + \frac{1}{\rho^2} R \frac{\partial^2 Q}{\partial \phi^2} = 0$$

$$\Rightarrow \frac{1}{R} \rho \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} R = - \frac{1}{Q} \frac{\partial^2 Q}{\partial \phi^2} = k^2$$

Clearly $Q(\phi) = \sin k\phi; \cos k\phi$ ← Since V has period 2π , k must be n , an integer

If we guess that $R(\rho) = \rho^l$ then

$$\rho \frac{\partial}{\partial \rho} \rho \frac{\partial}{\partial \rho} R = \rho \frac{\partial}{\partial \rho} \rho l \rho^{l-1} = \rho l^2 \rho^{l-1} = l^2 R$$

So this works if $l = \pm k = \pm n$

We conclude the solutions of $\nabla^2 V = 0$ have the form

$$V(\rho, \phi) = \rho^n \sin n\phi; \rho^n \cos n\phi$$

$$\rho^{-n} \sin n\phi; \rho^{-n} \cos n\phi$$

Clearly the top line are the ones to use for $\rho < a$ and the bottom line for $\rho > a$, to avoid divergences at $\rho = 0$ and $\rho = \infty$

2-2

$$\rho < a \quad V(\rho, \phi) = \sum_{n=1}^{\infty} \rho^n (a_n \cos n\phi + b_n \sin n\phi)$$

$$\rho > a \quad V(\rho, \phi) = \sum_{n=1}^{\infty} \rho^{-n} (c_n \cos n\phi + d_n \sin n\phi)$$

We want V continuous at $\rho = a$ and also that the discontinuity in $\frac{\partial V}{\partial \rho}$, the normal piece of \vec{E} be $\sigma/\epsilon_0 = \sigma_0/\epsilon_0 \cos 4\phi$.

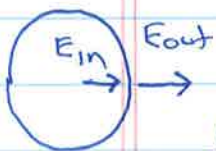
This suggests we try a solution with $b_n = d_n = 0$ and only a_4 and c_4 nonzero

$$\rho < a \quad V(\rho, \phi) = \rho^4 a_4 \cos 4\phi$$

$$\rho > a \quad V(\rho, \phi) = \rho^{-4} c_4 \cos 4\phi$$

$$\text{Continuity of } V(a, \phi) \Rightarrow a^4 a_4 = a^{-4} c_4 \Rightarrow c_4 = a^8 a_4$$

$$\begin{aligned} \text{Discontinuity in } E_n &\Rightarrow +4\rho^{-5} c_4 \cos 4\phi + 4\rho^3 a_4 \cos 4\phi \Big|_{\rho=a} \\ &= \frac{1}{\epsilon_0} \sigma_0 \cos 4\phi \\ &= -\frac{\partial V}{\partial \rho} \end{aligned}$$



Cancelling the $\cos 4\phi$ factors, and using $c_4 = a^8 a_4$

$$E_{out} A - E_{in} A = \sigma_0 A$$

$$\underbrace{\hspace{10em}}_{\phi \epsilon}$$

$$E_{out} - E_{in} = \sigma_0 / \epsilon_0$$

$$4a^{-5} a^8 a_4 + 4a^3 a_4 = \sigma_0 / \epsilon_0$$

$$a_4 = \sigma_0 / 8\epsilon_0 a^3$$

$$V(\rho, \phi) = \begin{cases} \sigma_0 / 8\epsilon_0 a^3 \rho^4 \cos 4\phi & \rho < a \\ \sigma_0 / 8\epsilon_0 a^3 \rho^{-4} \cos 4\phi & \rho > a \end{cases}$$

3-1

↙ Extra credit.

3. Same starting point as problem 2, except

$$V(\rho, \phi, z) = R(\rho)Q(\phi)T(z)$$

$$\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial V}{\partial r} + \frac{1}{r^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} = 0 \quad \leftarrow \text{Using Laplacian in cylindrical coordinates}$$

$$Q \cdot T \frac{d^2 R}{dr^2} + \frac{Q \cdot T}{r} \frac{dR}{dr} + \frac{R \cdot T}{r^2} \frac{d^2 Q}{d\phi^2} + R \cdot Q \frac{d^2 T}{dz^2} = 0$$

divide
by RQT

↪

$$\frac{1}{R} \frac{d^2 R}{dr^2} + \frac{1}{rR} \frac{dR}{dr} + \frac{1}{r^2 Q} \frac{d^2 Q}{d\phi^2} = - \frac{1}{T} \frac{d^2 T}{dz^2}$$

⏟

r, ϕ dependence
only

⏟

z dependence
only

$\therefore = \text{constant}$. Let's call it $-\ell^2$

$$T(z) = a e^{\ell z} + b e^{-\ell z} \quad \leftarrow \frac{d^2 T}{dz^2} = \ell^2 T$$

(or $\cosh \ell z, \sinh \ell z$)

$$\frac{r^2}{R} \frac{d^2 R}{dr^2} + \frac{r}{R} \frac{dR}{dr} + \ell^2 r^2 = - \frac{1}{Q} \frac{d^2 Q}{d\phi^2}$$

⏟

r only

⏟

ϕ only

$\therefore = \text{constant}$. Let's call it m^2

$$Q(\phi) = c e^{im\phi} + d e^{-im\phi} \quad \leftarrow \frac{d^2 Q}{d\phi^2} = -m^2 Q$$

(or $\sin m\phi, \cos m\phi$)

Note: m must be an integer since $V(\phi + 2\pi) = V(\phi)$

The r dependence is ugly!

$$r^2 \frac{d^2 R}{dr^2} + r \frac{dR}{dr} + (\ell^2 r^2 - m^2) = 0$$

Define $x = \ell r$

$$x^2 \frac{d^2 R}{dx^2} + x \frac{dR}{dx} + (x^2 - m^2) = 0$$

This is the Bessel Eqn \uparrow

You can stop here ... But I cite a few items

$$m \text{ noninteger} \left\{ \begin{array}{l} J_m(x) = x^m \sum_0^{\infty} \frac{(-1)^n x^{2n}}{2^{m+2n} n! \Gamma(m+n+1)} \\ J_{-m}(x) = x^{-m} \sum_0^{\infty} \frac{(-1)^n x^{2n}}{2^{-m+2n} n! \Gamma(-m+n+1)} \end{array} \right.$$

$$R(x) = e J_m(x) + f J_{-m}(x)$$

$$\Gamma(v) \equiv \int_0^{\infty} e^{-t} t^{v-1} dt$$

m integer $J_{-m}(x) = (-1)^m J_m(x)$ and instead second soln to Bessel eqn is

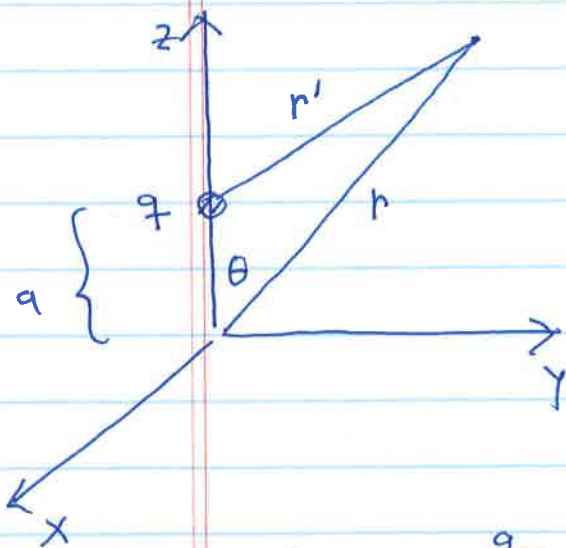
$$Y_m(x) = \lim_{p \rightarrow m} \frac{J_p(x) \cos p\pi - J_{-p}(x)}{\sin p\pi}$$

$$R(x) = e J_m(x) + f Y_m(x)$$

See posted notes for extensive discussion

and applications: Diffusion, Schroedinger, Stat Mech!

[4] We did this in class...



$$V = \frac{1}{4\pi\epsilon_0} \frac{q}{r'}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{q}{(r^2 - 2ra\cos\theta + a^2)^{1/2}}$$

$$V(r, \theta) = \frac{q}{4\pi\epsilon_0 r} \left\{ 1 - \frac{1}{2} \left(-2 \frac{a\cos\theta}{r} + \frac{a^2}{r^2} \right) + \left(-\frac{1}{2} \right) \left(-\frac{3}{2} \right) \left(\frac{1}{2} \right) \frac{(-2 \frac{a\cos\theta}{r} + \frac{a^2}{r^2})^2}{r^2} + \dots \right\}$$

$$= \frac{q}{4\pi\epsilon_0 r} \left\{ 1 + \frac{a}{r} \cos\theta + \frac{a^2}{r^2} \left(\frac{3}{2} \cos^2\theta - \frac{1}{2} \right) + \dots \right\}$$

$$\uparrow$$

$$P_0(\cos\theta)$$

monopole

$$\frac{q}{4\pi\epsilon_0 r}$$

$$\uparrow$$

$$P_1(\cos\theta)$$

dipole

$$\frac{aq}{4\pi\epsilon_0 r^2 \cos\theta}$$

$$\uparrow$$

$$P_2(\cos\theta)$$

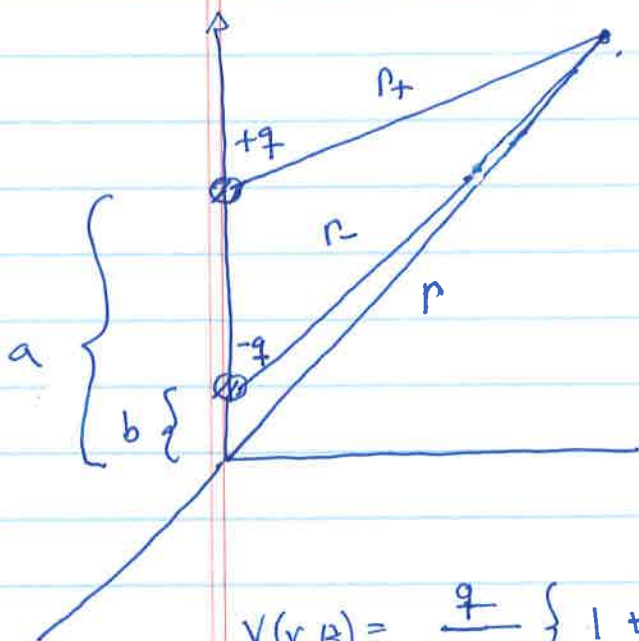
quadrupole

$$\frac{qa^2}{4\pi\epsilon_0 r^3} \left(\frac{3}{2} \cos^2\theta - \frac{1}{2} \right)$$

If \$a=0\$ only the monopole term survives.

5-1

5



$$V(r, \theta) = \frac{q}{4\pi\epsilon_0} \left\{ \frac{1}{r_+} - \frac{1}{r_-} \right\}$$

$$= \frac{q}{4\pi\epsilon_0} \left\{ \frac{1}{(r^2 - 2ra \cos\theta + a^2)^{1/2}} - \frac{1}{(r^2 - 2rb \cos\theta + b^2)^{1/2}} \right\}$$

Following #4 we see:

$$V(r, \theta) = \frac{q}{4\pi\epsilon_0} \left\{ 1 + \frac{a}{r} \cos\theta + \frac{a^2}{r^2} \left(\frac{3}{2} \cos^2\theta - \frac{1}{2} \right) + \dots \right. \\ \left. - 1 - \frac{b}{r} \cos\theta - \frac{b^2}{r^2} \left(\frac{3}{2} \cos^2\theta - \frac{1}{2} \right) + \dots \right\}$$

* There is no monopole term if $q_{\text{tot}} = +q - q = 0$.

* The dipole term is $\frac{q}{4\pi\epsilon_0} \frac{1}{r^2} (a-b) \cos\theta$

and depends only on the $+q/-q$ separation.

It is unaffected by shifting the charges by a common amount. Put alternately, the dipole term is invariant under changes of the origin.