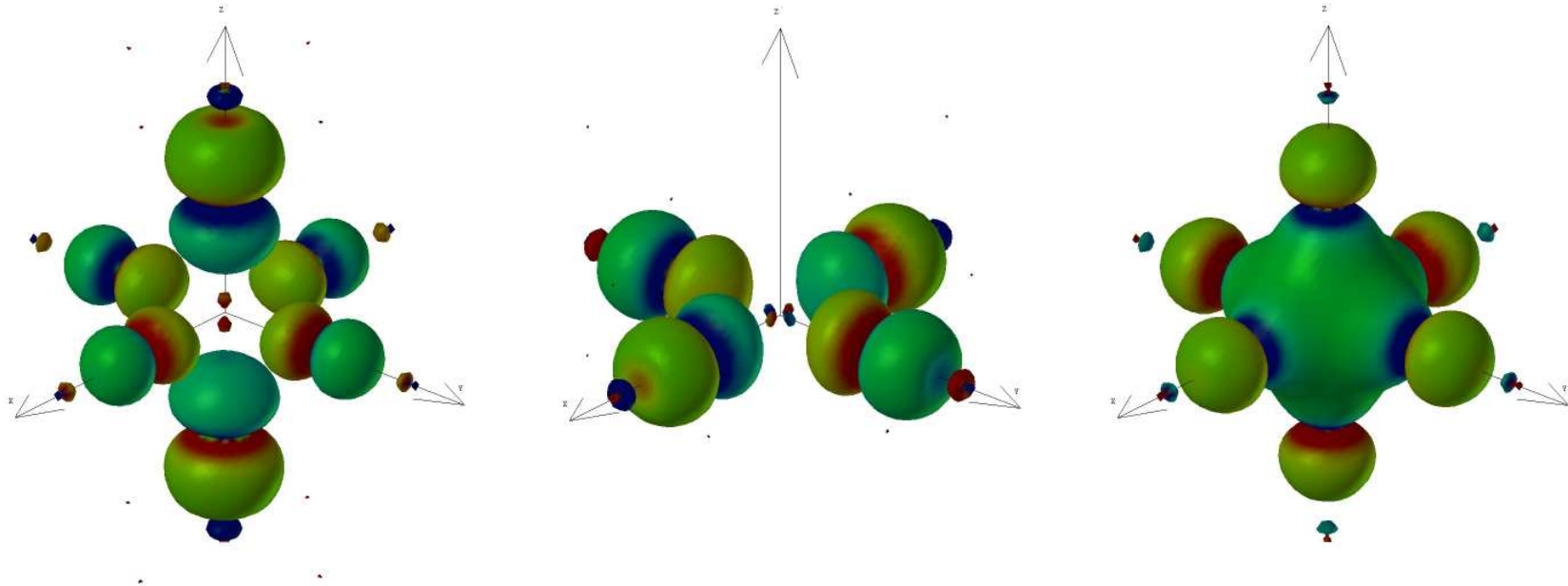


“Super atom” approach to local excitations in strongly correlated systems



Chi-Cheng Lee Brookhaven National Laboratory

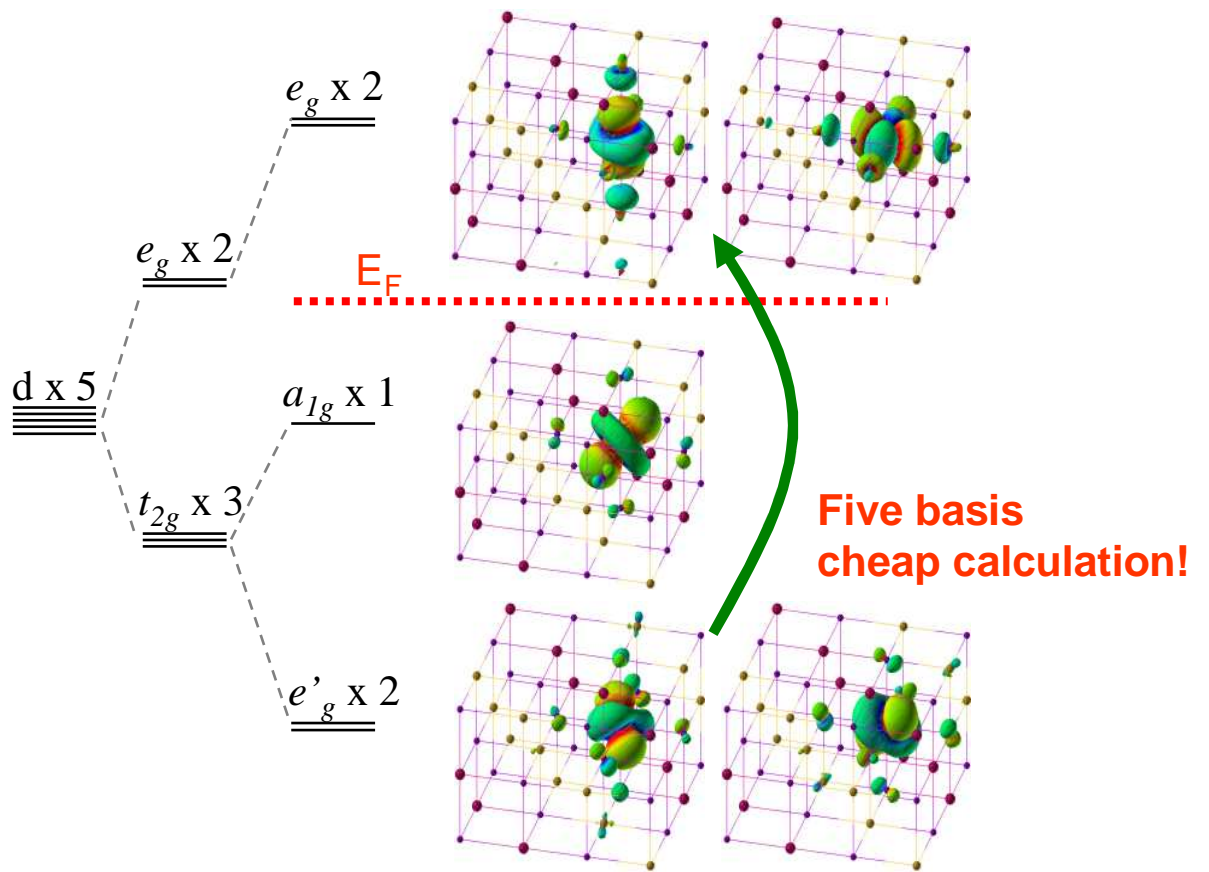
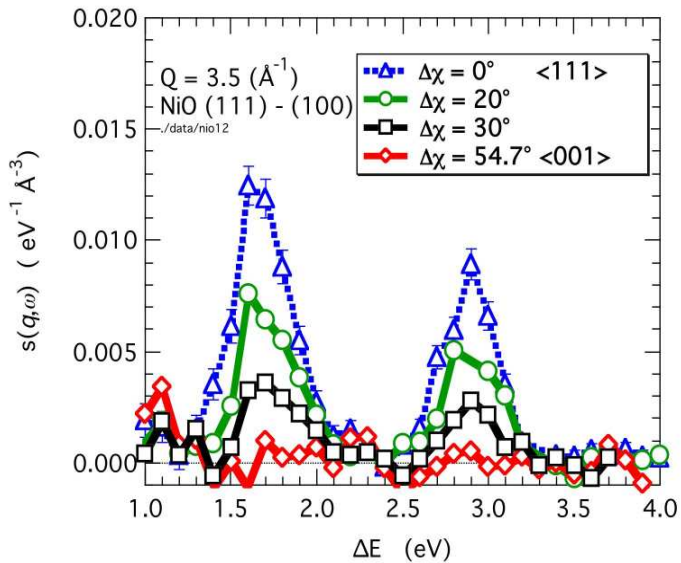
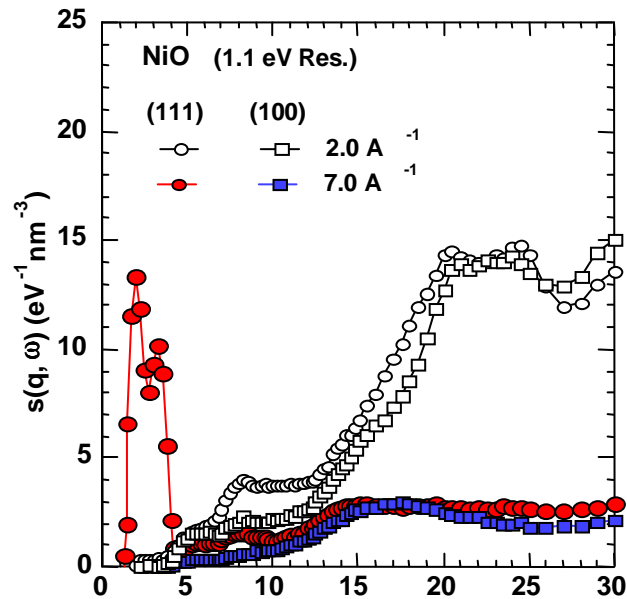
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SUNY Stony Brook University

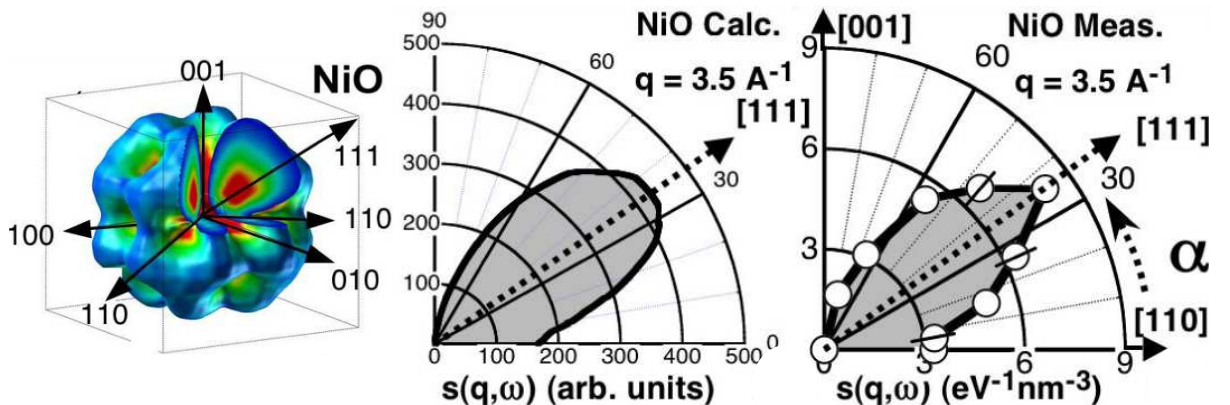
Introduction

- **NIXS found in-gap *d-d* excitation in Mott-insulator NiO**
- **BSE with Hartree-Fock-type kernel not working**
- **Local many-body Hamiltonian investigation**
- **Super atom concept in Wannier basis**
- **Many-body excitation and density response function**

NIXS experiment for NiO*



$$\chi(\mathbf{q}, \mathbf{q}'; \omega) = \sum_{11'} \langle 1 | e^{-i\mathbf{q} \cdot \hat{\mathbf{x}}} | 1' \rangle \langle 1' | e^{i\mathbf{q}' \cdot \hat{\mathbf{x}}} | 1 \rangle L(1, 1'; 1', 1; \omega)$$



*B. C. Larson *et al.*, PRL **99**, 026401 (2007)

Linear response within TDDFT via LDA+U energy functional (TDLDA+U)

equation of motion

$$G = G_0 + G_0 v_s G$$

$$v_s = v_{ext} + v_{Hartree} + v_{XC} + v_{local\ Hartree} - v_{local\ Fock} - [U(\text{loop}) - J(\text{loop})] \dots$$

response function

$$L = \frac{\delta G}{\delta v_{ext}} = -G \frac{\delta G^{-1}}{\delta v_{ext}} G$$

$$\chi = L + L \chi$$

(Bethe-Salpeter equation)

$$I = \text{Hartree} + f_{xc}^{LDA}(w) - (U-J) + \text{local Hartree} - \text{local Fock p-h attraction}$$

Hartree
long-range screening

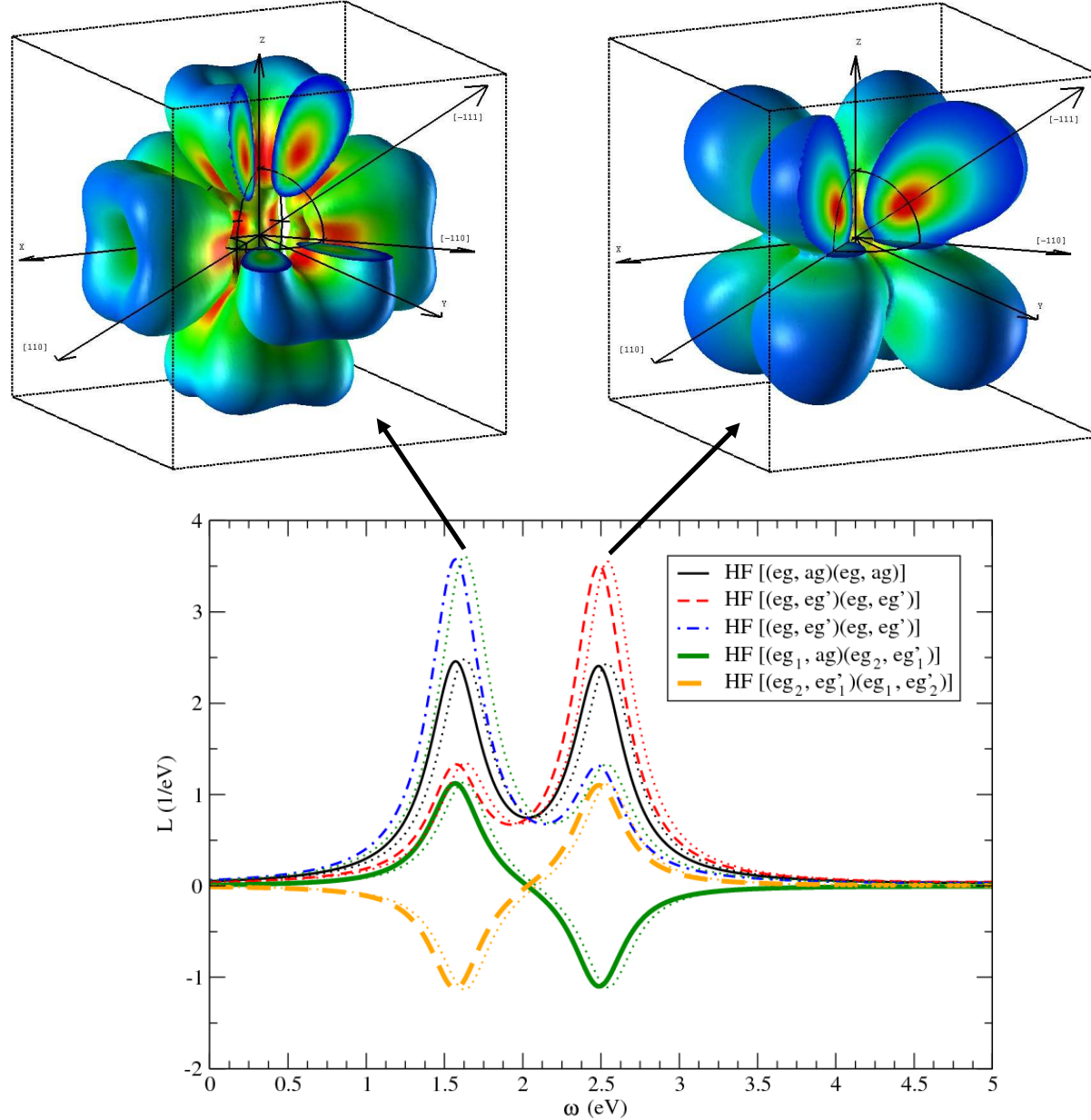
$f_{xc}^{LDA}(w)$

d.c.

local Hartree

local Fock
p-h attraction

Bethe-Salpeter solution with local Hartree-Fock interaction



Perturbation not working ! → Many-body diagonalization is necessary !

Local many-body Hamiltonian for strongly correlated systems

H_{LDA+U} :

$$E^{LDA+U}[\rho^\sigma(\vec{r}), \{n^\sigma\}] = E^{LSDA}[\rho^\sigma(\vec{r})] + E^U[\{n^\sigma\}] - E_{dc}[\{n^\sigma\}]$$

$$E^U[\{n\}] = \frac{1}{2} \sum_{\{m\}, \sigma} \{ \langle m, m'' | V_{ee} | m', m''' \rangle n_{mm'}^\sigma n_{m''m'''}^{-\sigma} + (\langle m, m'' | V_{ee} | m', m''' \rangle - \langle m, m'' | V_{ee} | m''', m' \rangle) n_{mm'}^\sigma n_{m''m'''}^\sigma \}$$

$$V_{mm'}^\sigma = \sum_{m''m'''} \{ \langle m', m'' | V_{ee} | m, m''' \rangle n_{m''m'''}^{-\sigma} + (\langle m', m'' | V_{ee} | m, m''' \rangle - \langle m', m'' | V_{ee} | m''', m \rangle) n_{m''m'''}^\sigma \}$$

$$-U(n - \frac{1}{2}) + J(n^\sigma - \frac{1}{2})$$

$$H_{LDA+U} = \sum_{Rm, R'm', \sigma} t_{Rm, R'm'}^\sigma C_{Rm\sigma}^+ C_{R'm'\sigma} = \sum_{Rm, R'm', \sigma} \langle WF_{Rm} | \hat{H}_{LDA+U}^\sigma | WF_{R'm'} \rangle C_{Rm\sigma}^+ C_{R'm'\sigma}$$

$$H = H_{local} + H_{nonlocal} = T_{ij}^\sigma C_{i\sigma}^+ C_{j\sigma} + \frac{1}{2} V_{ijkl} C_{i\sigma}^+ C_{j\sigma}^+ C_{l\sigma} C_{k\sigma} + T_{ij}^\sigma C_{i\sigma}^+ C_{j\sigma}$$

H_{SCHF} :

$$H_{local}^{HF} = T_{ij}^\sigma C_{i\sigma}^+ C_{j\sigma} + (V_{kij}^{eff} \langle C_{k\sigma}^+ C_{l\sigma} \rangle - \delta_{\sigma\sigma'} V_{kijl}^{eff} \langle C_{k\sigma}^+ C_{l\sigma} \rangle) C_{i\sigma}^+ C_{j\sigma}$$

$$H_{local}^{HF} = T_{ij}^\sigma C_{i\sigma}^+ C_{j\sigma} + \varepsilon^{-1} U_{ij}^\sigma C_{i\sigma}^+ C_{j\sigma}; U_{ij}^\sigma = (V_{kijl}^{eff} - \delta_{\sigma\sigma'} V_{kijl}^{eff}) \langle C_{k\sigma}^+ C_{l\sigma} \rangle \leftarrow LDA+U$$

Fitting :

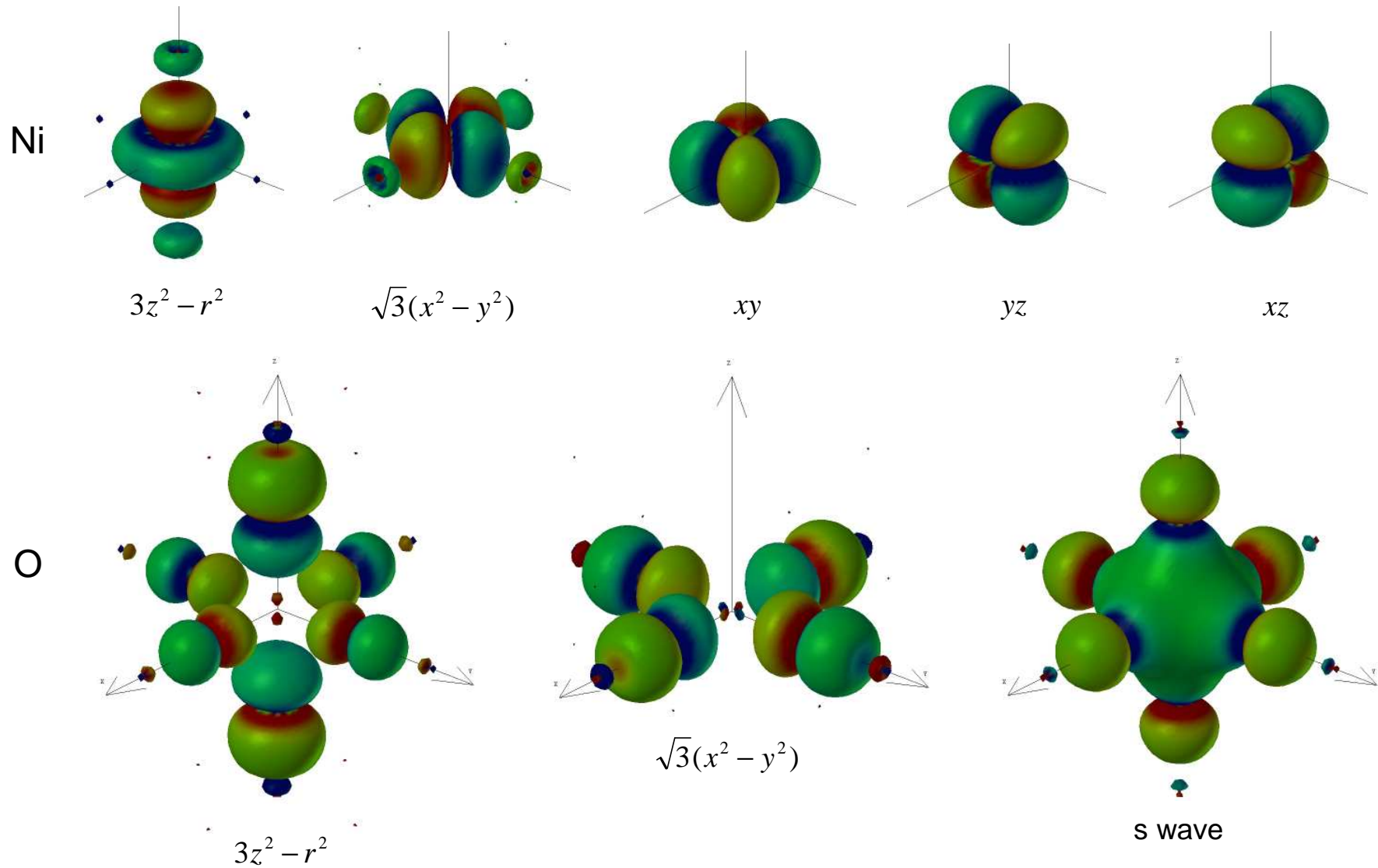
$$F = \sum_{ij\sigma} (t_{ij}^\sigma - T_{ij}^\sigma - \varepsilon^{-1} U_{ij}^\sigma)^2; T_{ii}^\uparrow = T_{ii}^\downarrow$$

Dielectric constant ~ 1

We can safely choose atomic Coulomb matrix elements in many-body Hamiltonian



Super atom for charge transfer insulator



Enlarge the contribution of “local atom”

$$H = H_{local} + H_{nonlocal}$$

(exact) (perturbation)

$$\langle WF_{0m} | \hat{H}_{LDA+U}^\sigma | WF_{0m'} \rangle :$$

		O									
		$3z^2 - r^2$	$\sqrt{3}(x^2 - y^2)$	yz	xz	xy	s	$3z^2 - r^2$	$\sqrt{3}(x^2 - y^2)$		
O	$3z^2 - r^2$	4.06	0.00	0.00	0.00	0.00	0.00	1.84	0.00	(up)	
	$\sqrt{3}(x^2 - y^2)$	0.00	4.06	0.00	0.00	0.00	0.00	0.00	1.84		→ t_{pd}
	yz	0.00	0.00	-2.95	0.00	0.00	0.00	0.00	0.00		
	xz	0.00	0.00	0.00	-2.95	0.00	0.00	0.00	0.00		
	xy	0.00	0.00	0.00	0.00	-2.95	0.00	0.00	0.00		
	s	0.00	0.00	0.00	0.00	0.00	-4.01	0.00	0.00		
	$3z^2 - r^2$	1.84	0.00	0.00	0.00	0.00	0.00	-2.02	0.00		
	$\sqrt{3}(x^2 - y^2)$	0.00	1.84	0.00	0.00	0.00	0.00	0.00	-2.02		
O	$3z^2 - r^2$	-5.06	0.00	0.00	0.00	0.00	0.00	1.81	0.00	(down)	
	$\sqrt{3}(x^2 - y^2)$	0.00	-5.06	0.00	0.00	0.00	0.00	0.00	1.81		
	yz	0.00	0.00	-3.69	0.00	0.00	0.00	0.00	0.00		
	xz	0.00	0.00	0.00	-3.69	0.00	0.00	0.00	0.00		
	xy	0.00	0.00	0.00	0.00	-3.69	0.00	0.00	0.00		
	s	0.00	0.00	0.00	0.00	0.00	-4.06	0.00	0.00		
	$3z^2 - r^2$	1.81	0.00	0.00	0.00	0.00	0.00	-2.02	0.00		
	$\sqrt{3}(x^2 - y^2)$	0.00	1.81	0.00	0.00	0.00	0.00	0.00	-2.02		

Local density-density response function

$$H_{local} = T_{ij}^{\sigma} C_{i\sigma}^+ C_{j\sigma} + \frac{1}{2} V_{ijkl} C_{i\sigma}^+ C_{j\sigma}^+ C_{l\sigma} C_{k\sigma} \quad (\text{in Wannier basis})$$

one-particle Green's function
two-particle Green's function

$$\chi(\bar{q}, \bar{q}; \omega) = \sum_m \left(\frac{|\langle \Psi_m | \hat{\rho}(-\bar{q}) | \Psi_0 \rangle|^2}{\hbar\omega - (E_m - E_0) + i\eta} - \frac{|\langle \Psi_m | \hat{\rho}(\bar{q}) | \Psi_0 \rangle|^2}{\hbar\omega - (E_0 - E_m) - i\eta} \right)$$

O

T_{ij}^{local}

	$3z^2 - r^2$	$\sqrt{3}(x^2 - y^2)$	yz	xz	xy	s	$3z^2 - r^2$	$\sqrt{3}(x^2 - y^2)$
eg	0.00	0.00	0.00	0.00	0.00	0.00	1.84	0.00
0.00	eg	0.00	0.00	0.00	0.00	0.00	0.00	1.84
0.00	0.00	t2g↑	0.00	0.00	0.00	0.00	0.00	0.00
0.00	0.00	0.00	t2g↓	0.00	0.00	0.00	0.00	0.00
0.00	0.00	0.00	0.00	t2g↑	0.00	0.00	0.00	0.00
0.00	0.00	0.00	0.00	0.00	0.00	-4.03↓	0.00	0.00
1.84	0.00	0.00	0.00	0.00	0.00	0.00	-2.02↓	0.00
0.00	1.84	0.00	0.00	0.00	0.00	0.00	0.00	-2.02↓

(up)

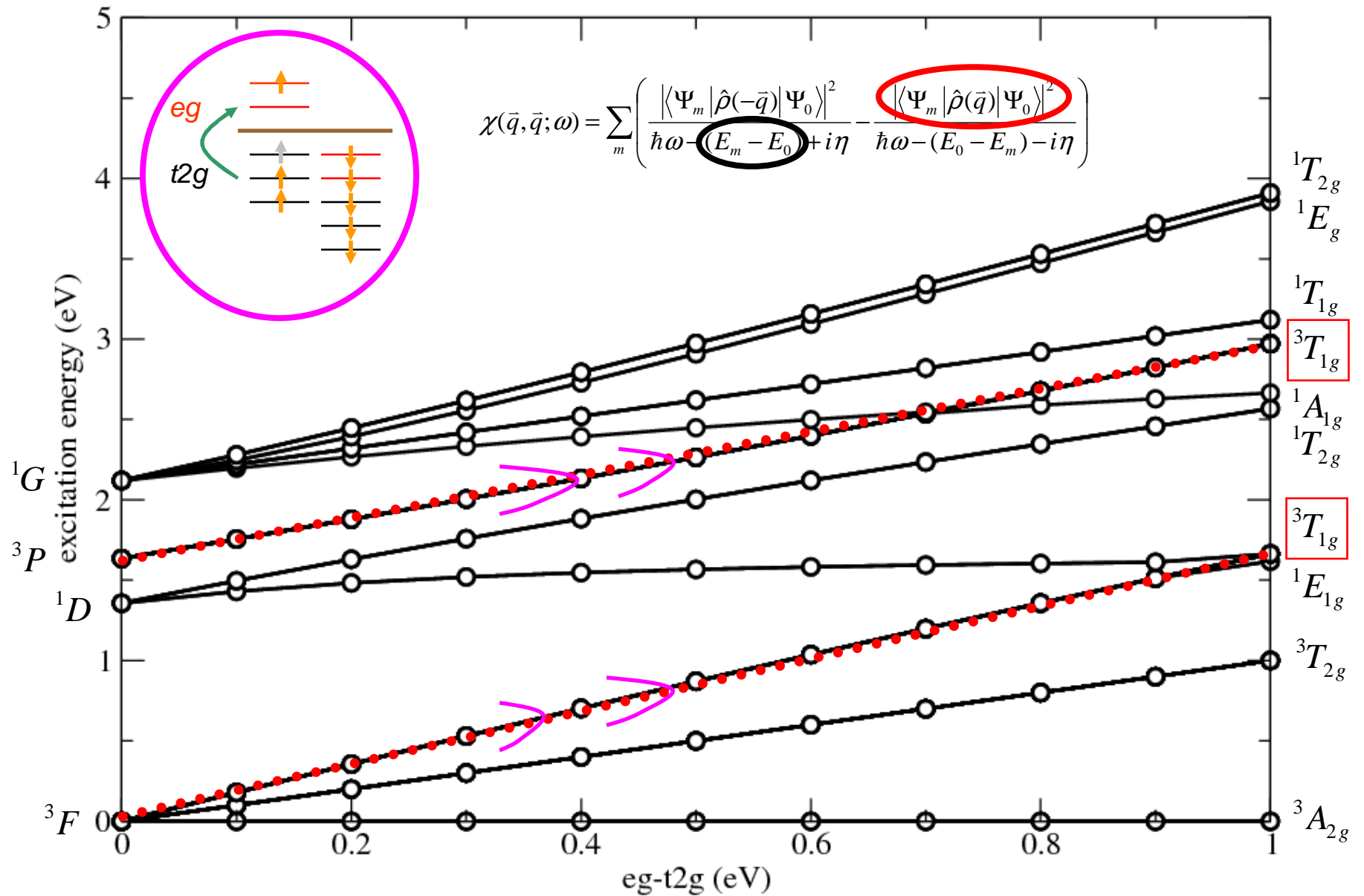
eg	0.00	0.00	0.00	0.00	0.00	0.00	1.81	0.00
0.00	eg	0.00	0.00	0.00	0.00	0.00	0.00	1.81
0.00	0.00	t2g	0.00	0.00	0.00	0.00	0.00	0.00
0.00	0.00	0.00	t2g	0.00	0.00	0.00	0.00	0.00
0.00	0.00	0.00	0.00	t2g	0.00	0.00	0.00	0.00
0.00	0.00	0.00	0.00	0.00	0.00	-4.03	0.00	0.00
1.81	0.00	0.00	0.00	0.00	0.00	0.00	-2.02	0.00
0.00	1.81	0.00	0.00	0.00	0.00	0.00	0.00	-2.02

(down)

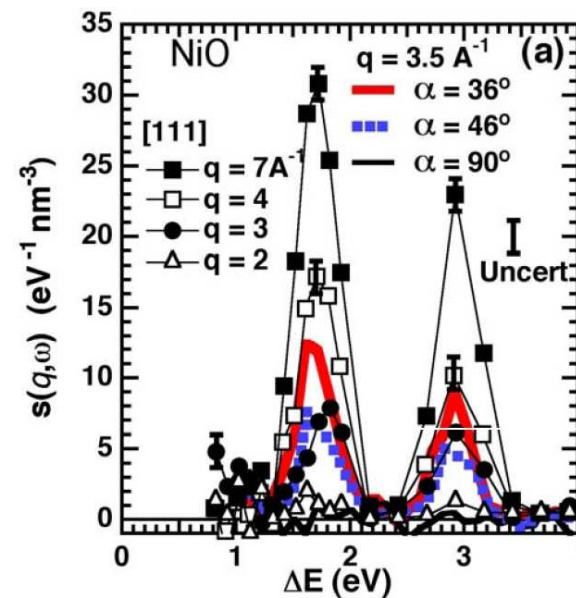
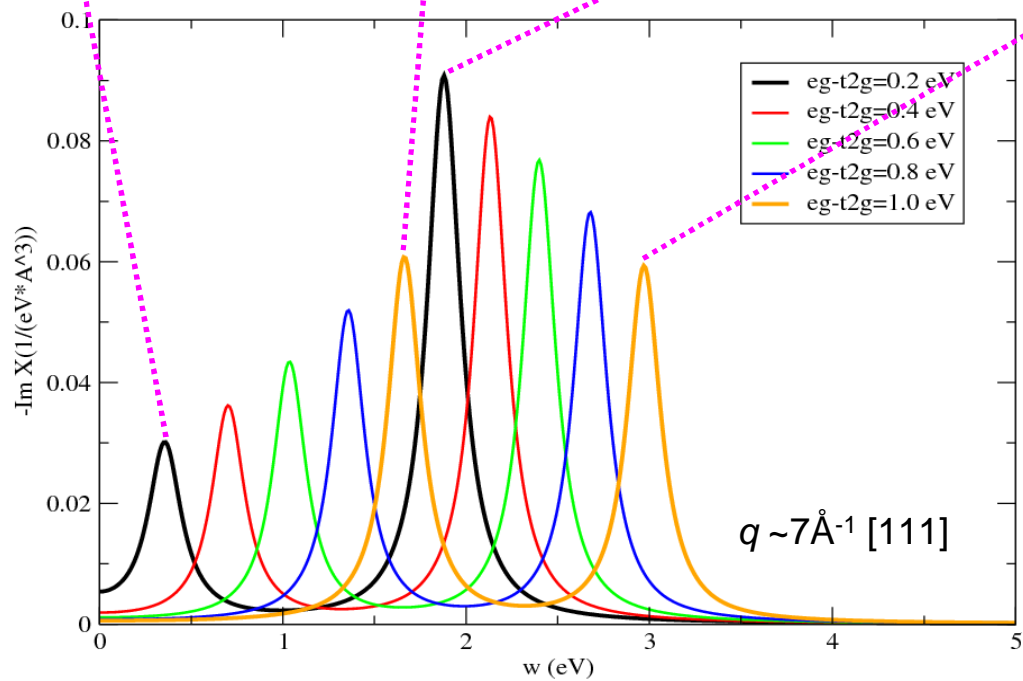
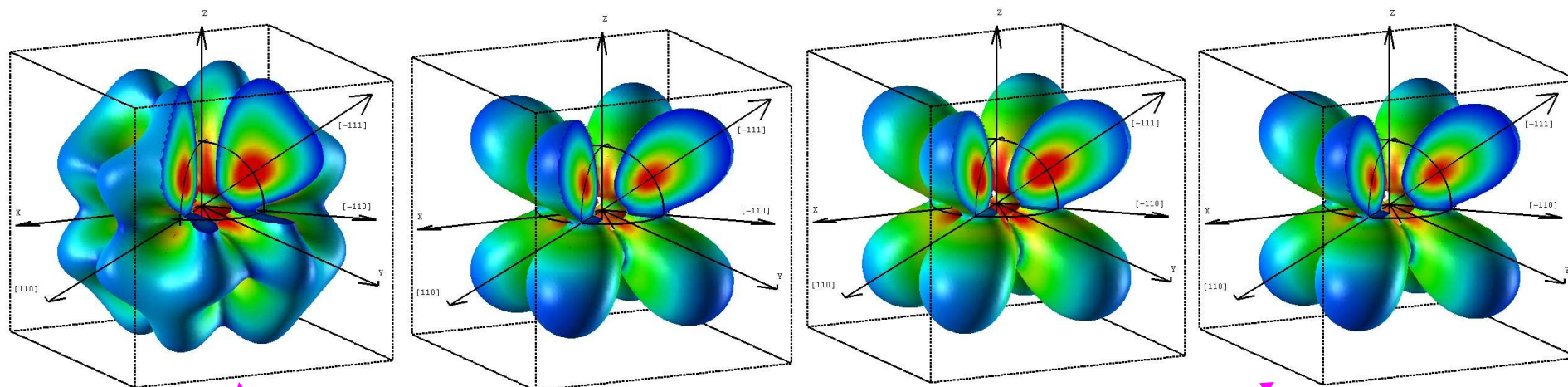
Many-body excitation for local atom

(eg-t2g=1eV to get experiment energy)

d orbitals of Ni only, no super atom of O included

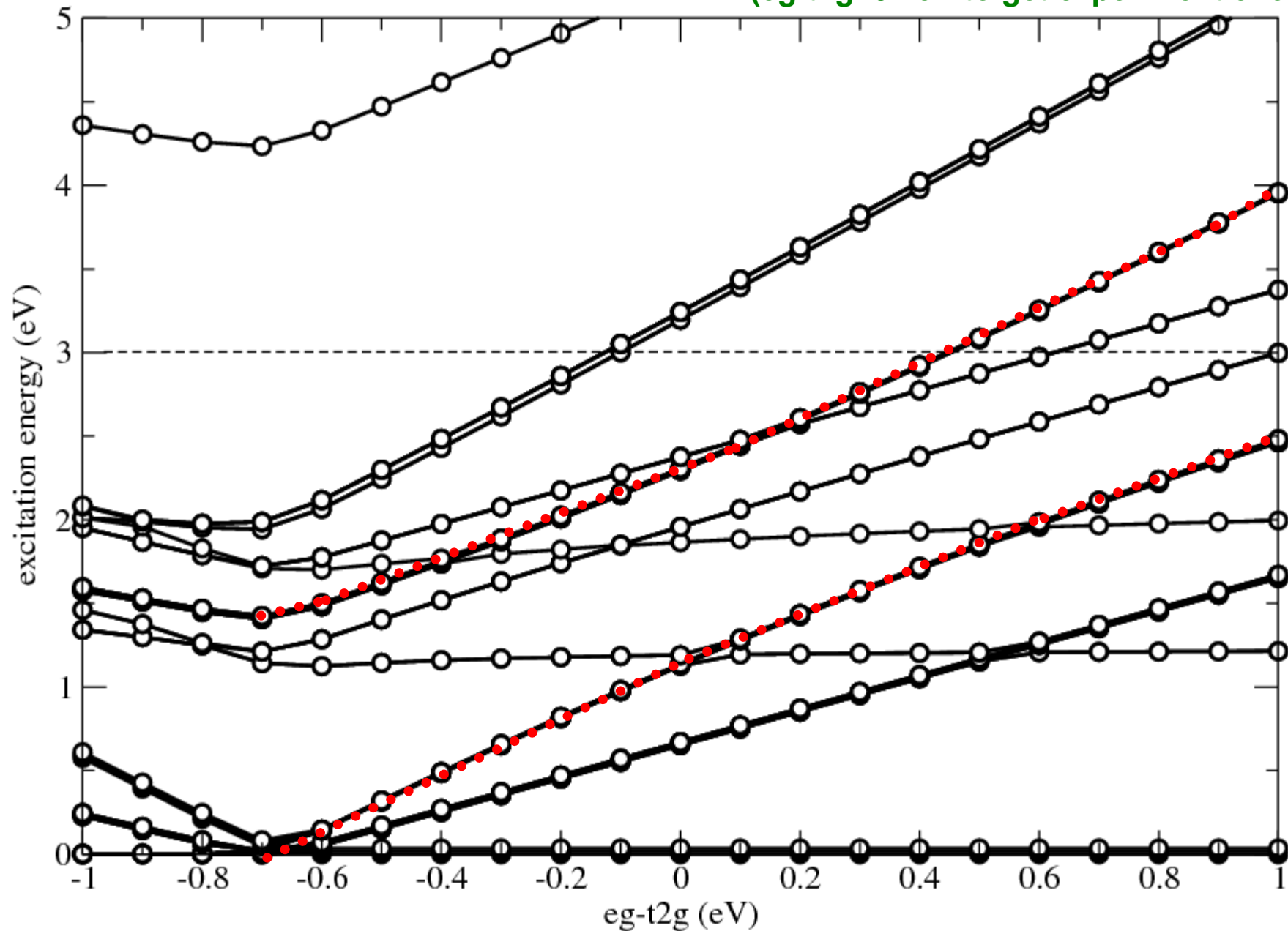


Density response for Ni only

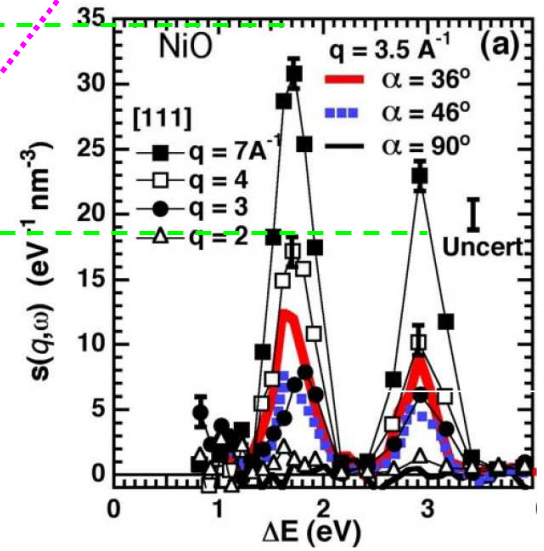
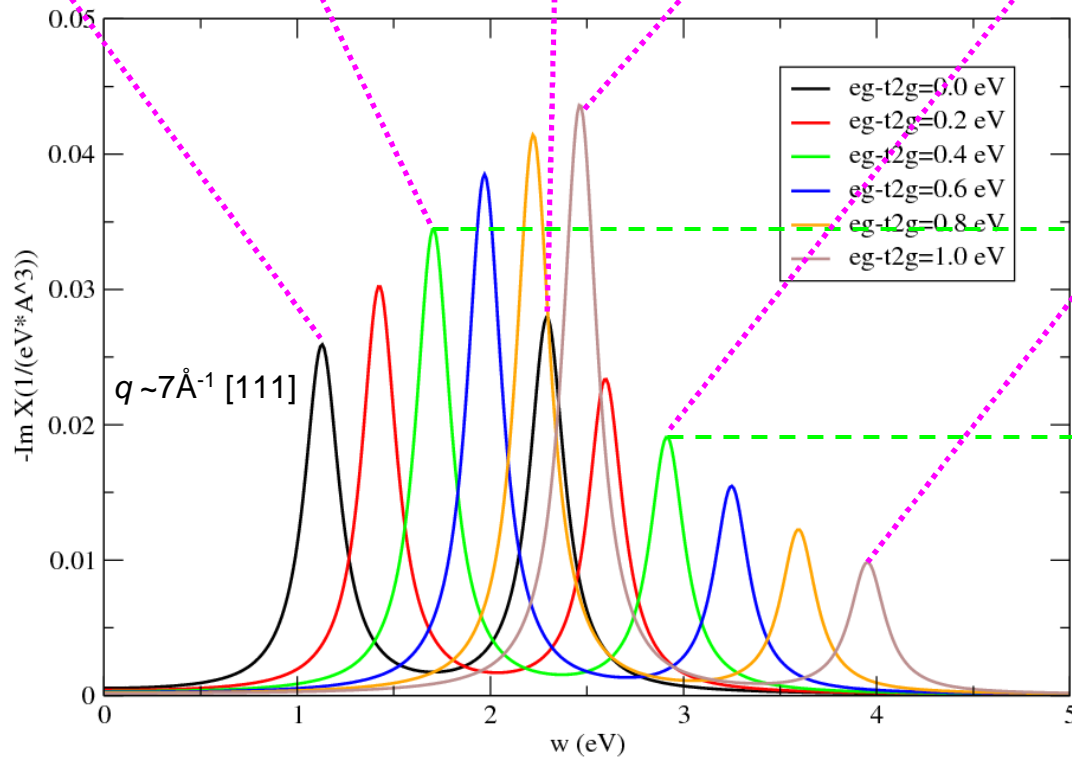
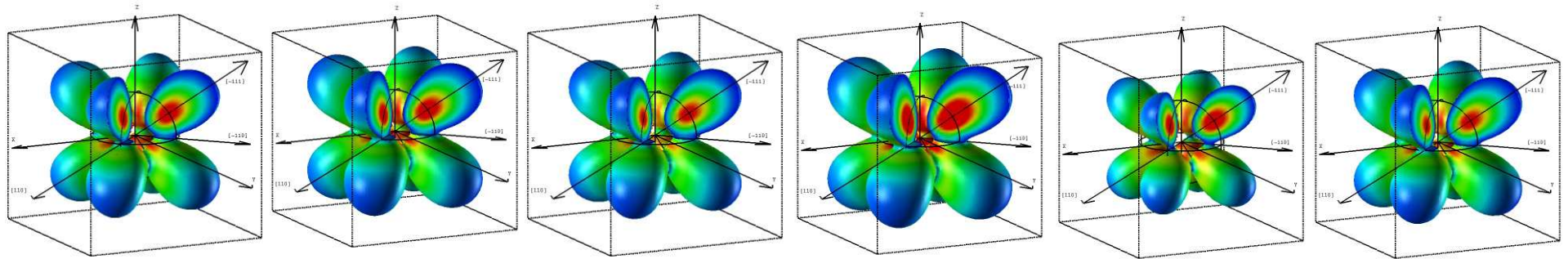


Many-body excitation with super atom of O

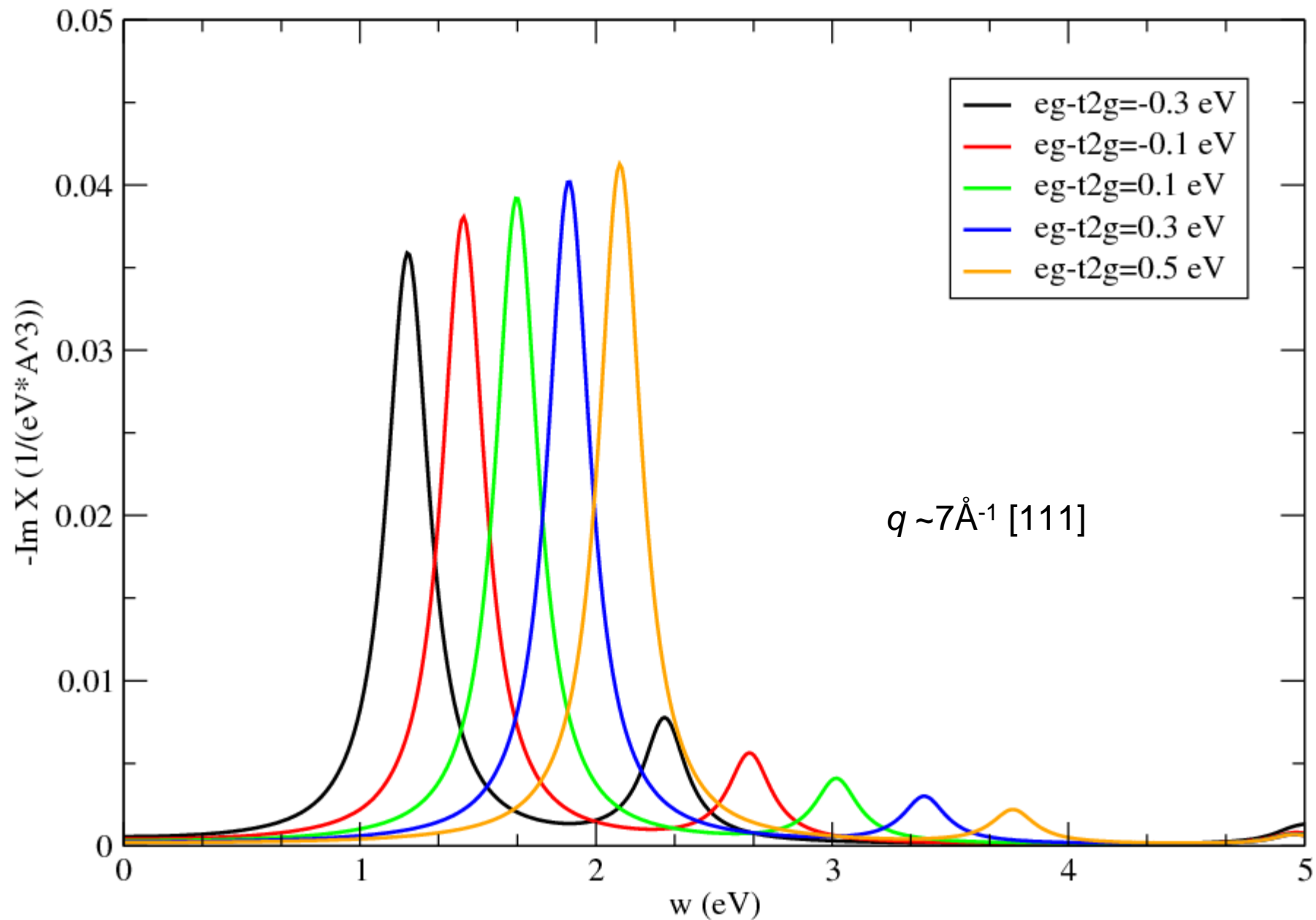
($eg-t2g=0.4\text{eV}$ to get experiment energy)



Density response with super atom

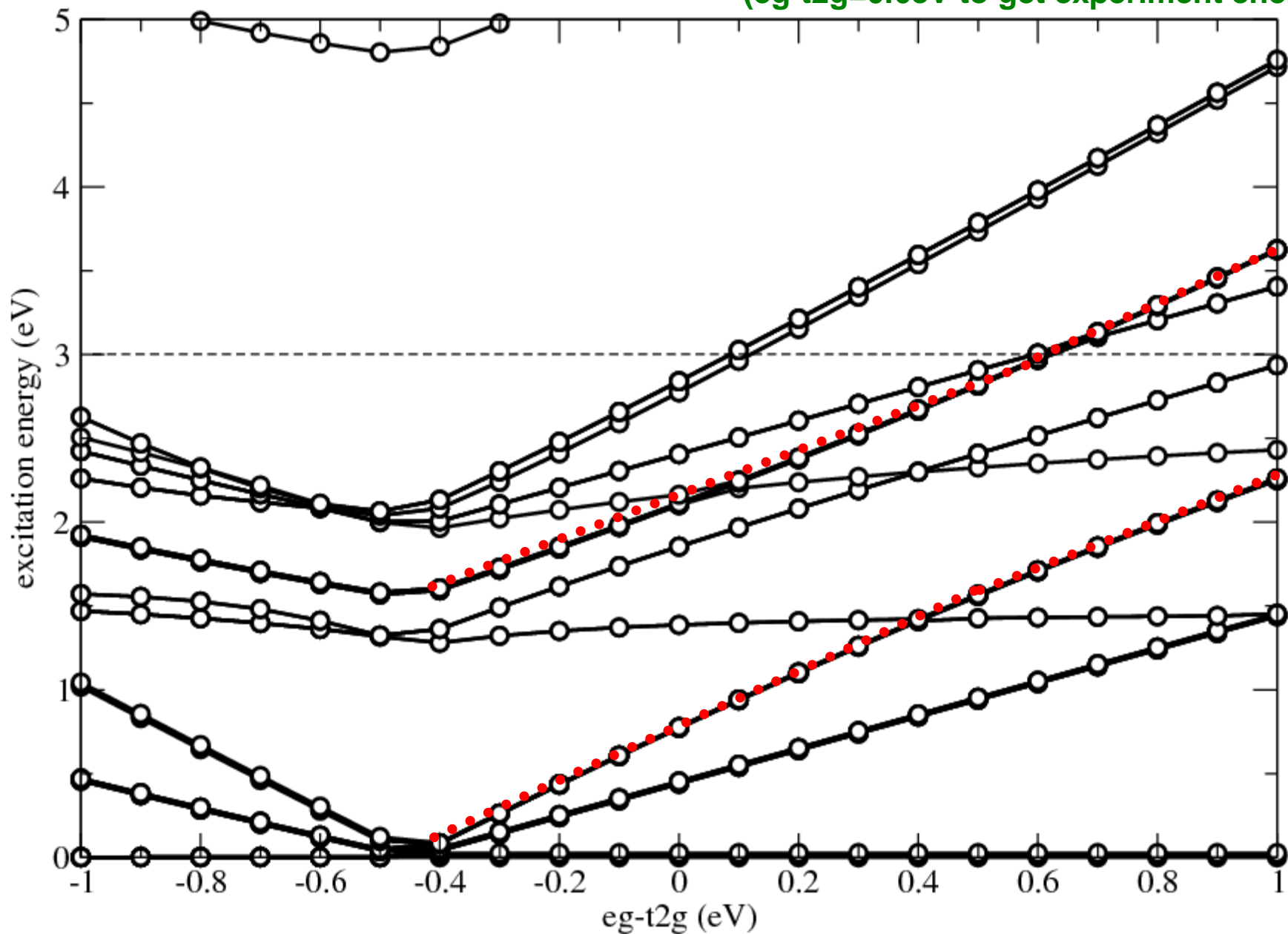


Density response with super atom (heighten O_{site} 3 eV)

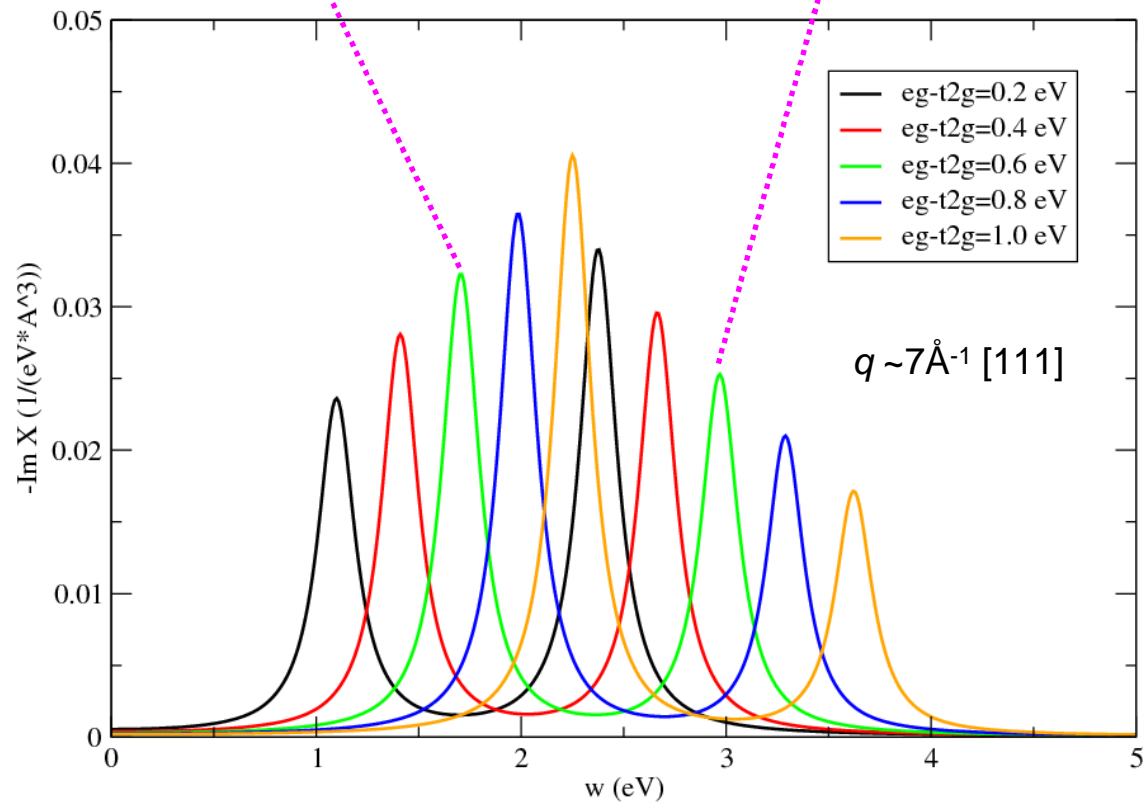
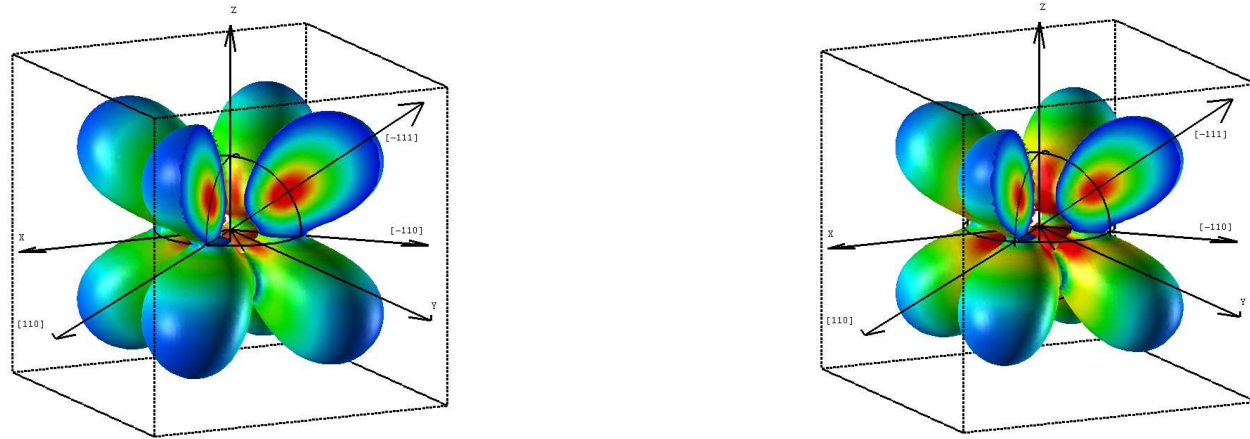


Many-body excitation with super atom of O (lower O_{site} 3 eV)

(eg-t2g=0.6eV to get experiment energy)



Density response with super atom (lower O_{site} 3 eV)

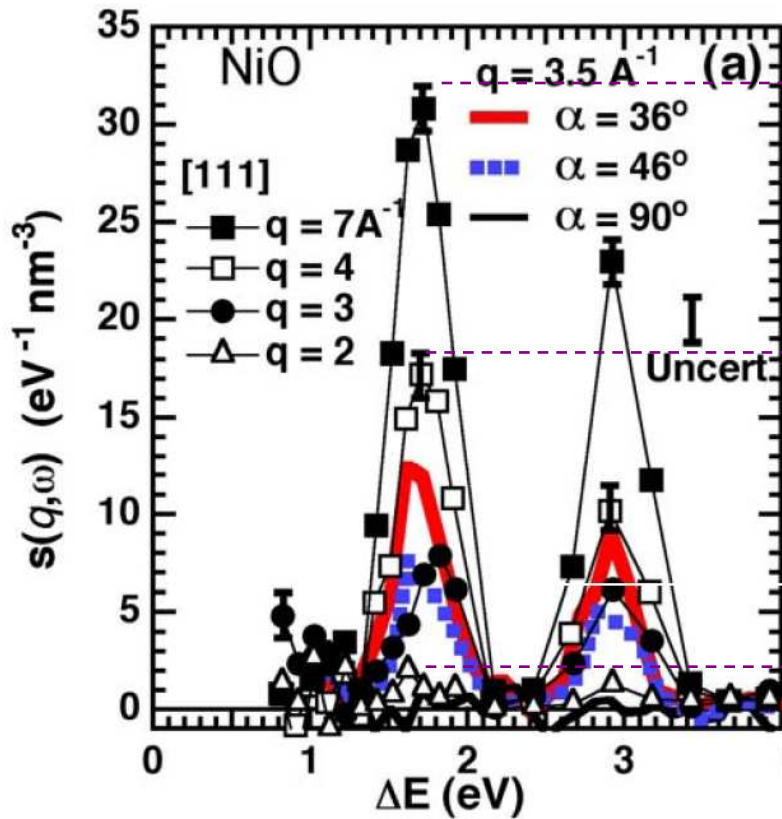


Density response with super atom

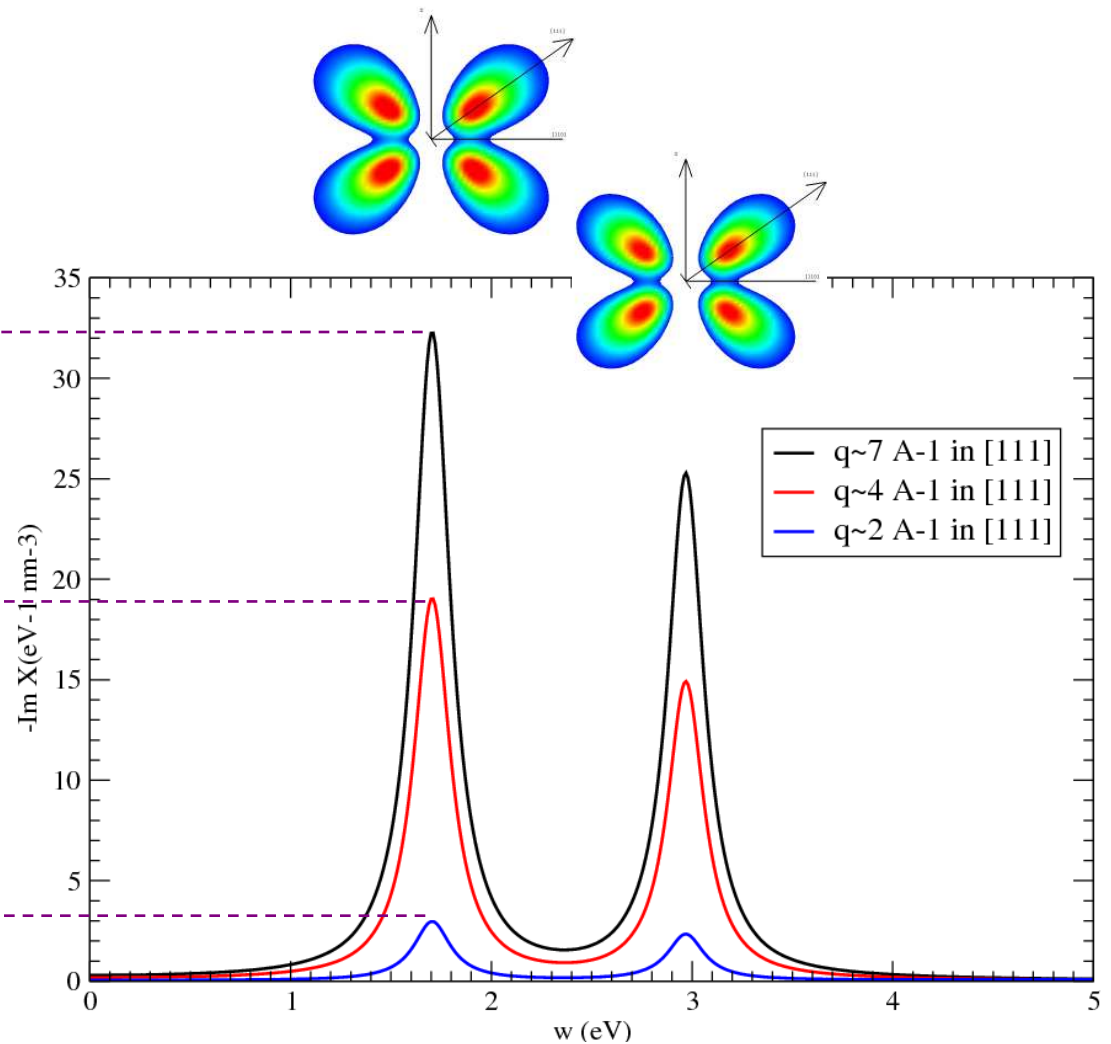
$Ni_{eg}-Ni_{t2g}=0.6\text{ eV}$

$Ni_{eg}-O_{eg}=-54.62\text{ eV}$

$V_{mm'm''m'''}: U=8\text{ eV}, J=0.95\text{ eV}$



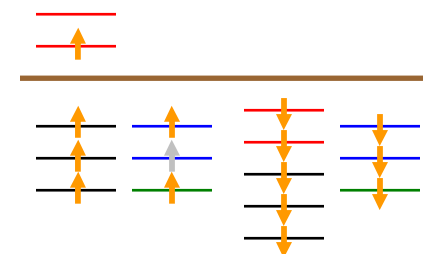
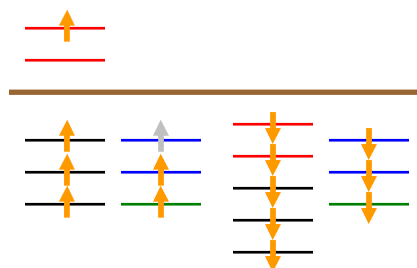
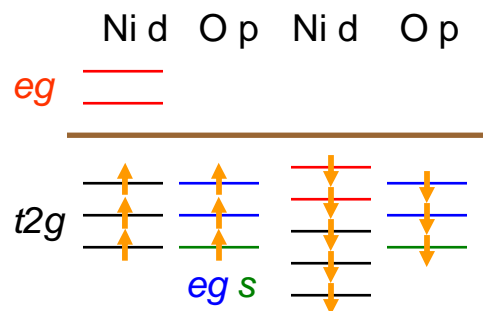
experiment



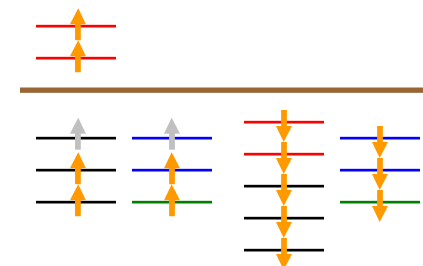
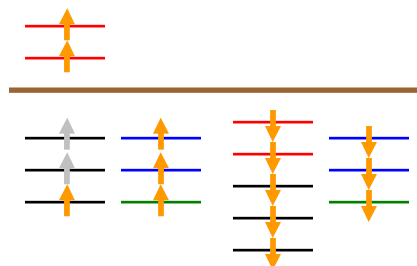
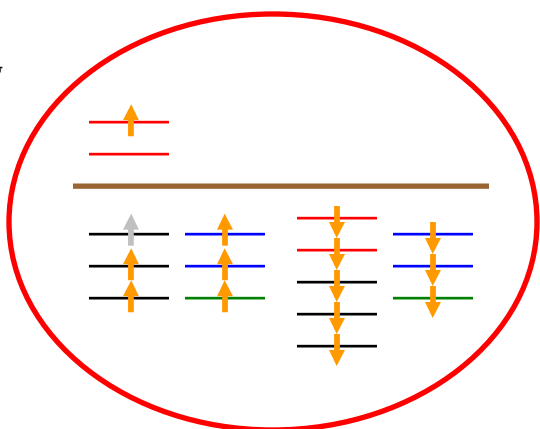
calculation

Many-body basis

GS



T_{1g}



measured by NIXS

Summary

- Treat local Coulomb interaction as perturbation → BSE → not working
- Diagonalizing local Hamiltonian composed of super atom
- Traditionally $10Dq(=e_g-t_{2g})=1\text{eV}$ provides correct excitation energy and q -dependence but not spectral weight
- Including oxygen changes the $10Dq$ to 0.6eV also help the spectral weight
- Many-body calculation of local Hamiltonian suggests same q -dependence for both peaks due to the contribution comes from the same many-body basis (one electron of t_{2g} excites to one e_g state)
- More possibility in investigating local many-body Hamiltonian e.g. quasiparticle calculation compared to ARPES

Next step

- First-principles site energies (e_g , t_{2g})
- Inclusion of non-local effect to study self energy → low-energy Hamiltonian